

Direct and indirect constraints on CP-violation in the Higgs sector

Jordy de Vries: Forschungszentrum Jülich (IKP-IAS)

Mostly based on: Arxiv:1510:00725

With: V. Cirigliano, Y-T. Chien, W. Dekens, E. Mereghetti
(Los Alamos National Laboratory)



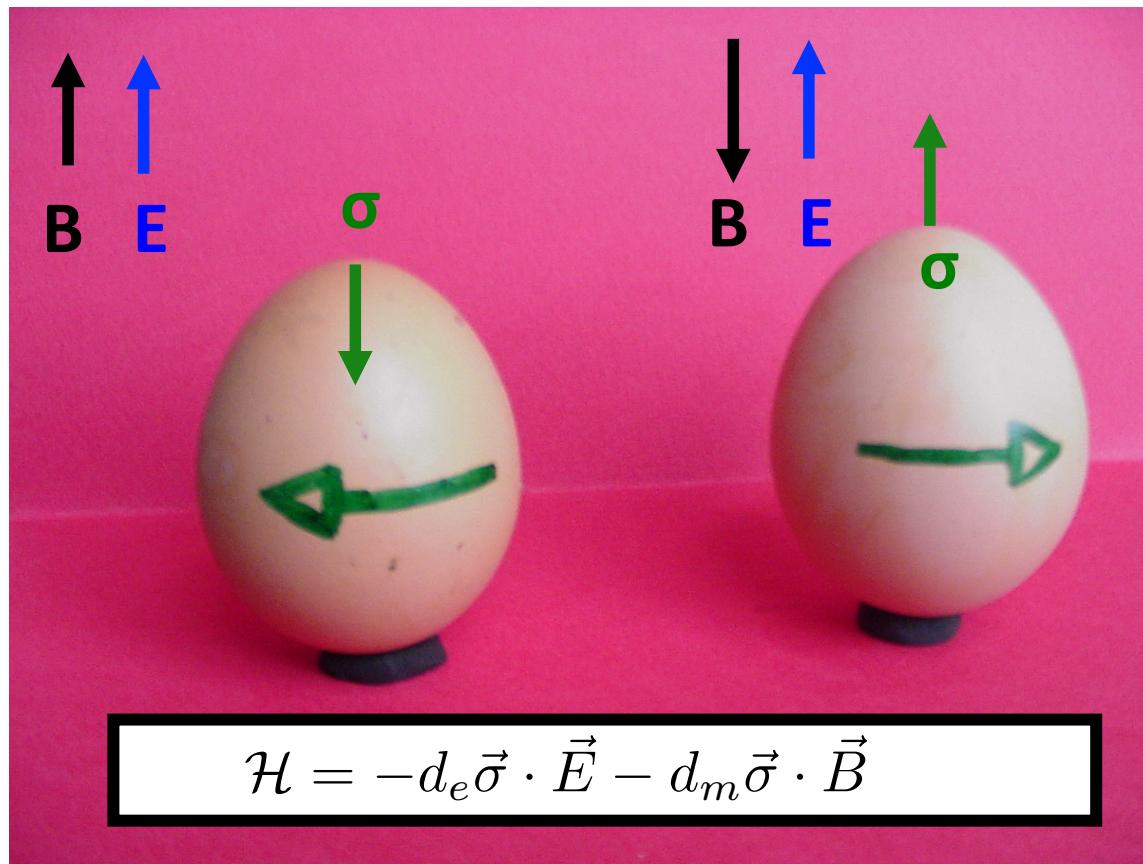
Outline of this talk

- **Part I:** The search for CP violation with EDMs.
- **Part II:** The EFT framework and its connection to low energy
 - EDMs of hadrons, nuclei, atoms, and molecules
- **Part III:** EDM/LHC Constraints on CPV Higgs couplings

EDMs 101

- Electric and Magnetic Dipole Moment (EDM and MDM)

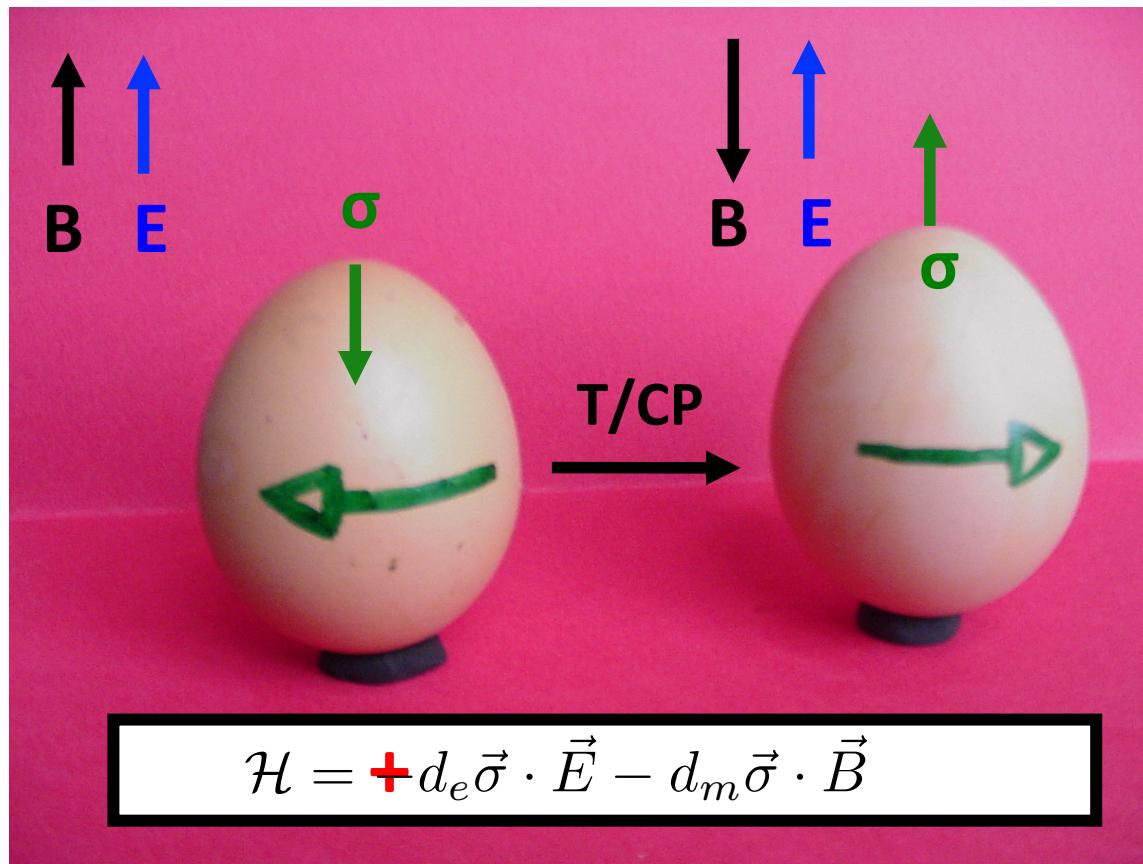
$$\mathcal{L}_d = -\frac{d_e}{2} \bar{\Psi} \sigma^{\mu\nu} \gamma^5 \Psi F_{\mu\nu} - \frac{d_m}{2} \bar{\Psi} \sigma^{\mu\nu} \Psi F_{\mu\nu}$$



EDMs 101

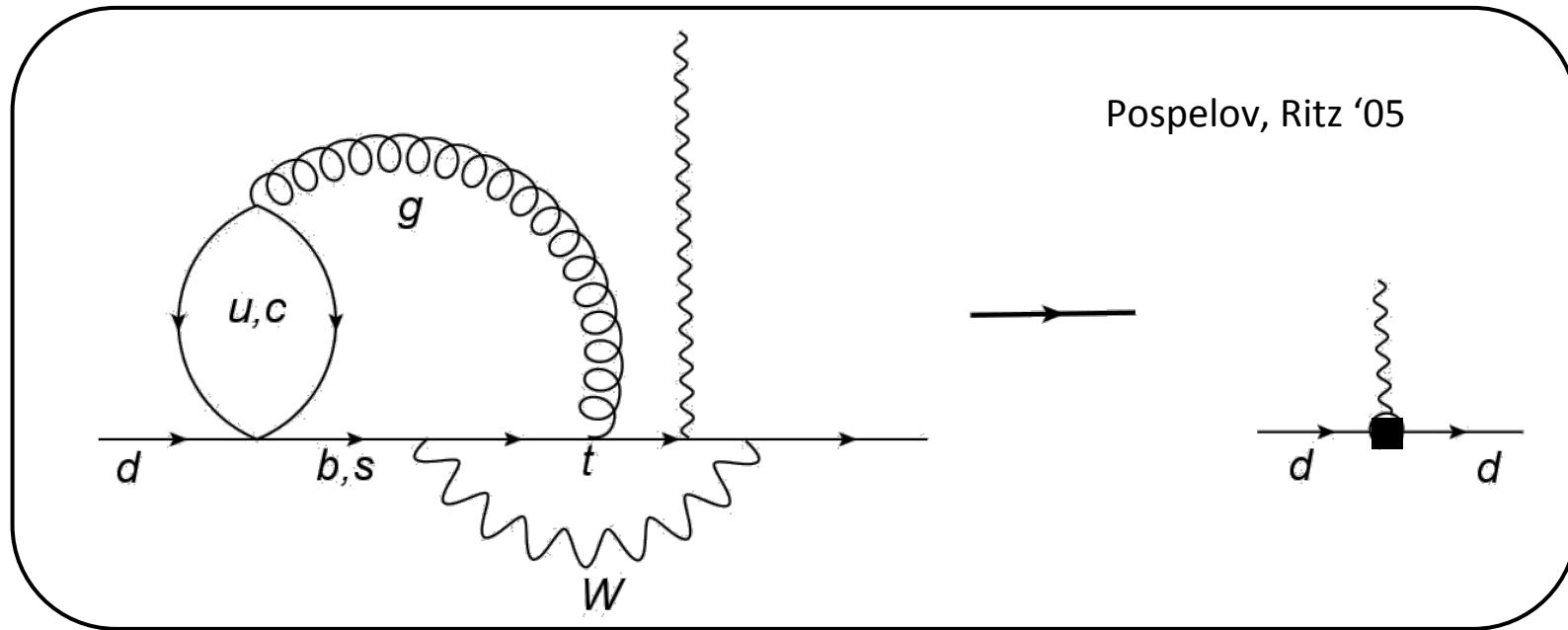
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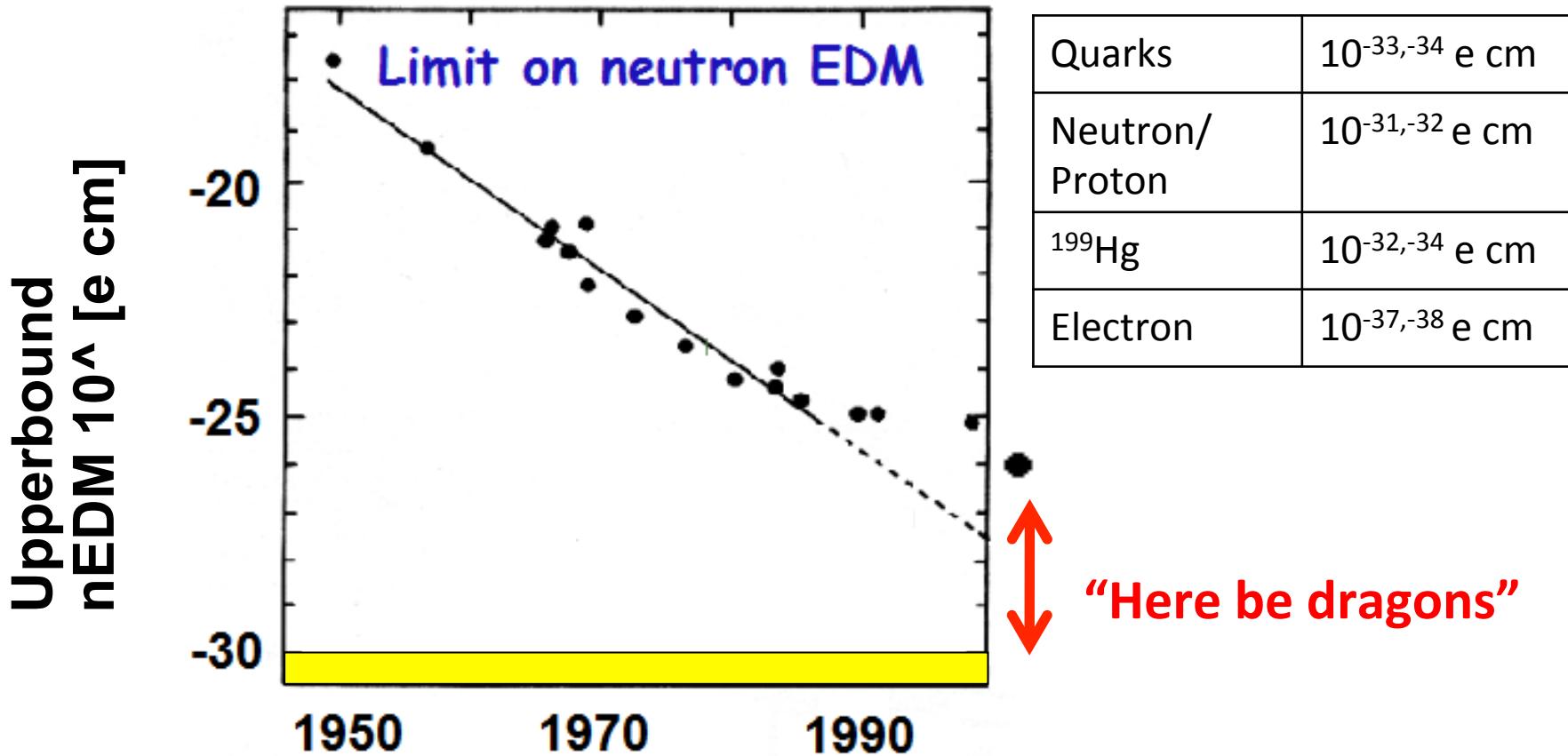
EDMs in the Standard Model

- Electroweak CP-violation very ineffective



- Quark EDMs = 0 at 2-loops , Electron EDM = 0 at 3-loops
- Dominant neutron EDM from CP-odd four-quark operators

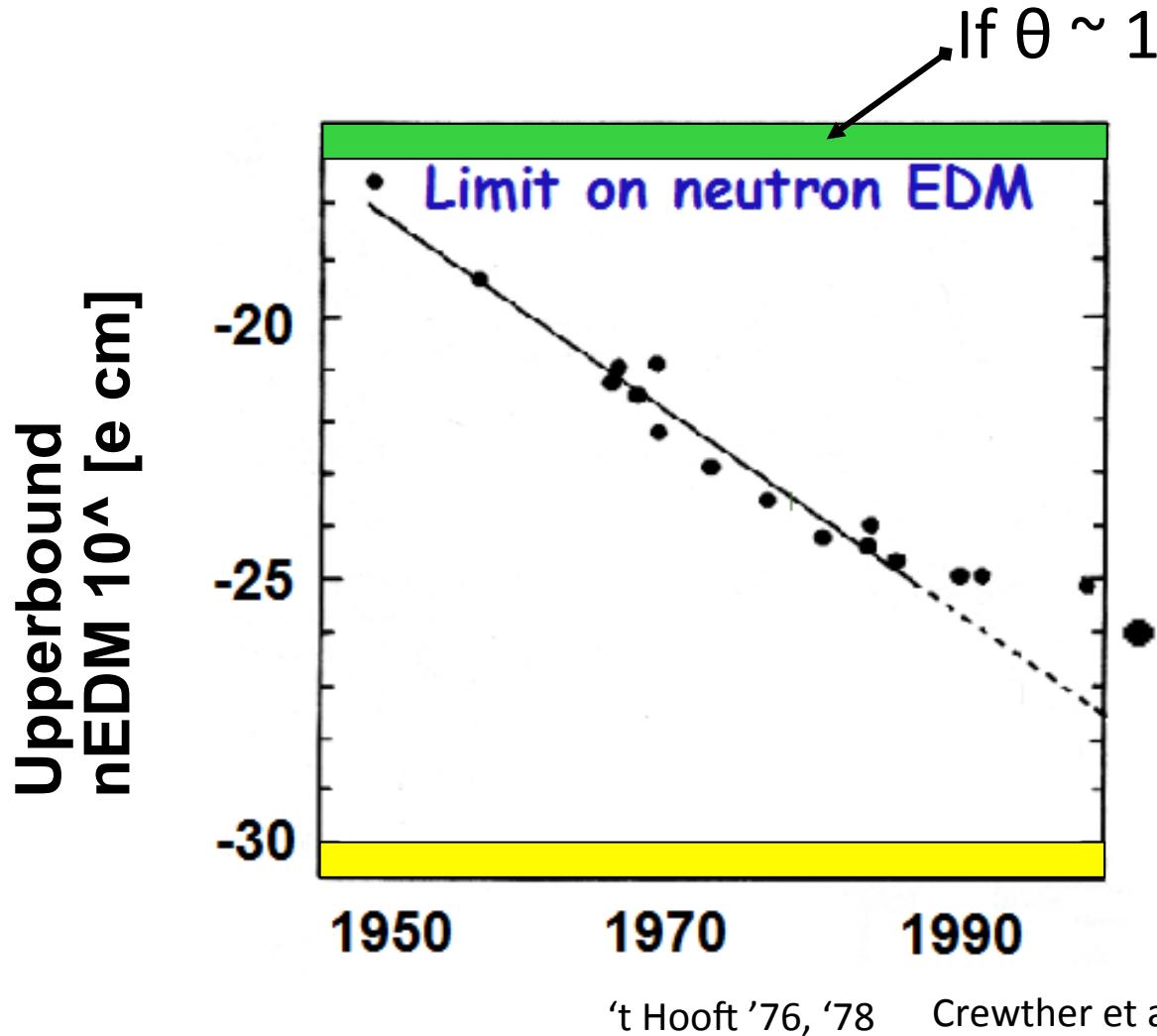
Neutron EDM from CKM



5 to 6 orders **below** upper bound \longleftrightarrow Out of reach!

With linear extrapolation: CKM neutron EDM in 2075....

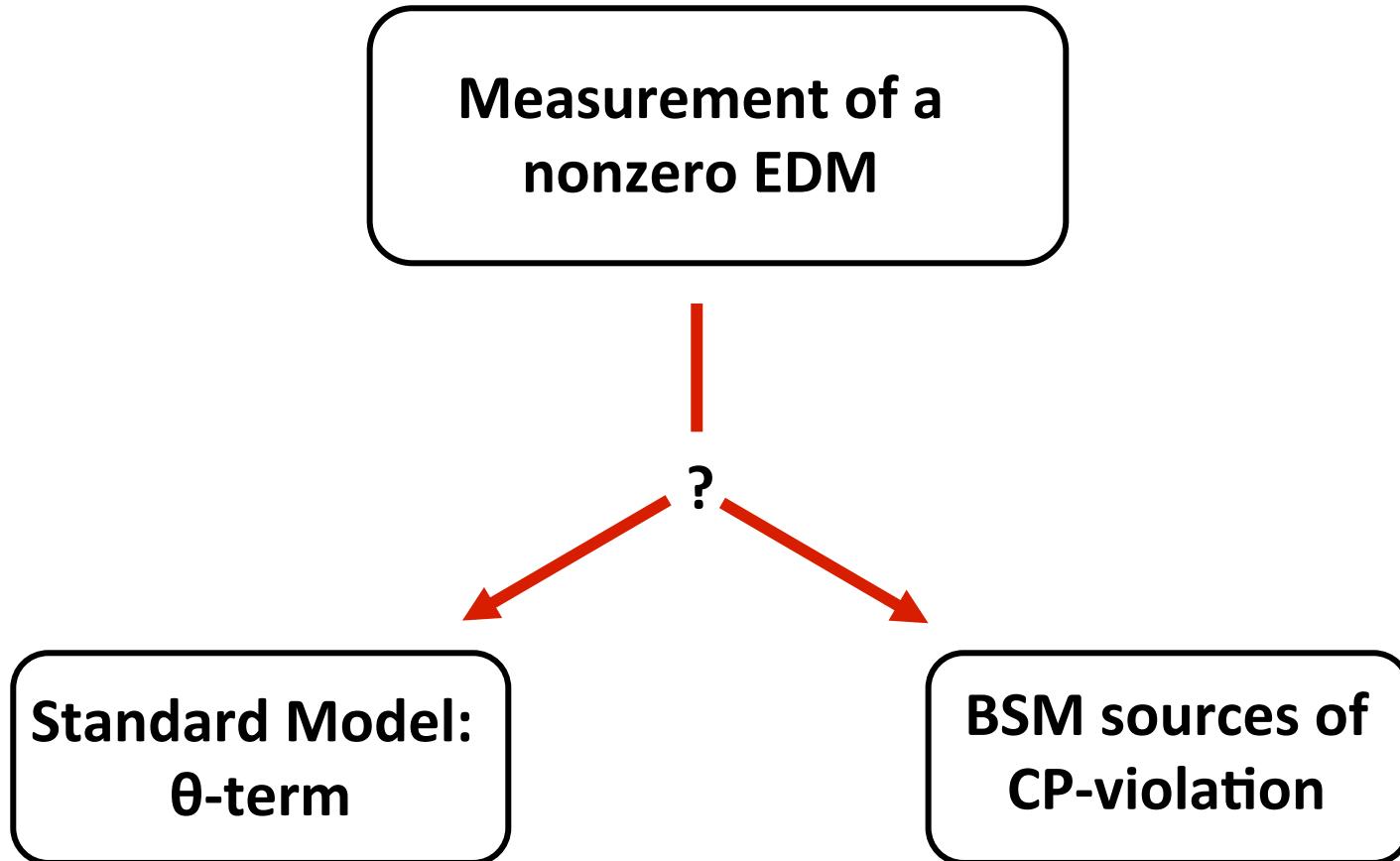
Neutron EDM from theta term



$$\theta \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu} G_{\alpha\beta}$$

Sets θ upper bound: $\theta < 10^{-10}$

In upcoming experiments:



For the foreseeable future: EDMs are
'background-free' searches for new physics

Active experimental field

System	Group	Limit	C.L.	Value	Year
^{205}TI	Berkeley	1.6×10^{-27}	90%	$6.9(7.4) \times 10^{-28}$	2002
YbF	Imperial	10.5×10^{-28}	90	$-2.4(5.7)(1.5) \times 10^{-28}$	2011
$\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$	Yale	6.05×10^{-25}	90	$-1.07(3.06)(1.74) \times 10^{-25}$	2012
PbO	Yale	1.7×10^{-26}	90	$-4.4(9.5)(1.8) \times 10^{-27}$	2013
ThO	ACME	8.7×10^{-29}	90	$-2.1(3.7)(2.5) \times 10^{-29}$	2014
n	Sussex-RAL-ILL	2.9×10^{-26}	90	$0.2(1.5)(0.7) \times 10^{-26}$	2006
^{129}Xe	UMich	6.6×10^{-27}	95	$0.7(3.3)(0.1) \times 10^{-27}$	2001
^{199}Hg	UWash	3.1×10^{-29}	95	$0.49(1.29)(0.76) \times 10^{-29}$	2009
muon	E821 BNL $g-2$	1.8×10^{-19}	95	$0.0(0.2)(0.9) \times 10^{-19}$	2009

$$d_e \leq 10^{-28} e \text{ cm} \simeq 10^{-14} e \text{ GeV}^{-1}$$

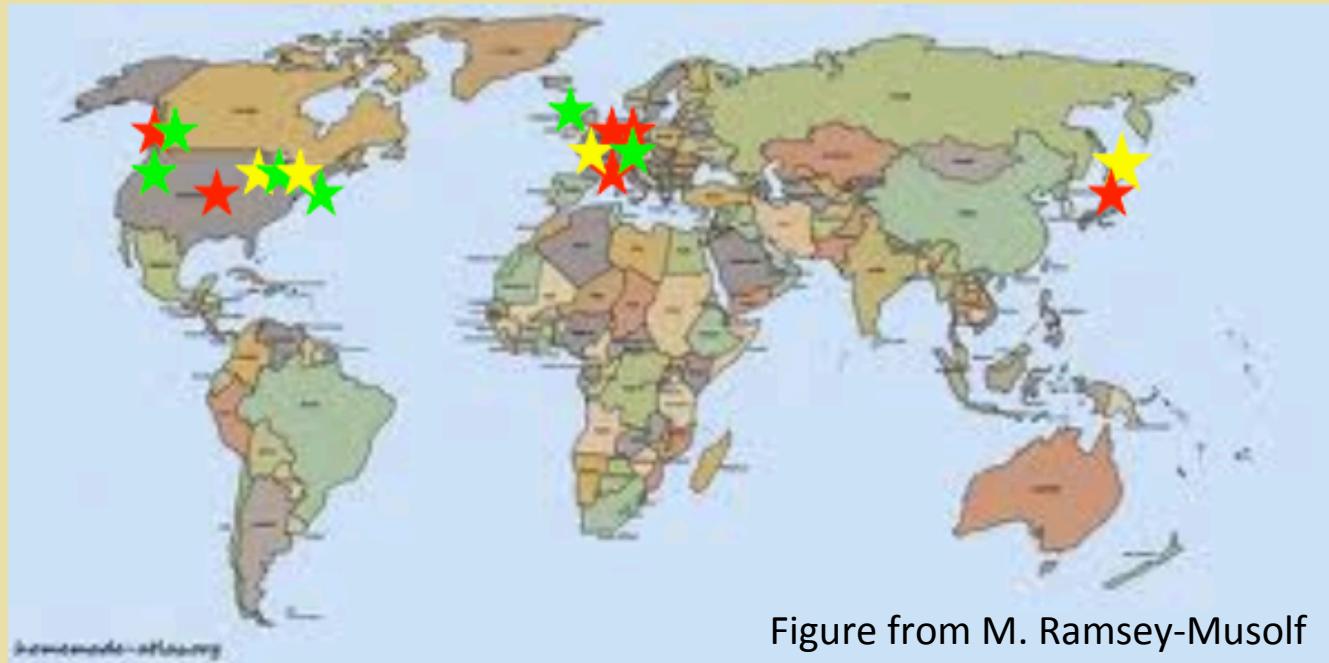
$$d_e \sim \left(\frac{\alpha_{em}}{\pi}\right)^n \frac{m_e}{\Lambda^2} \sin \phi$$

If phase = O(1): $\Lambda > 10 \text{ TeV}$ (n=1), $\Lambda > 0.5 \text{ TeV}$ (n=2)

(Model dependent!)

Active experimental field

System	Group	Limit	C.L.	Value	Year
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~ 100 x better sensitivity

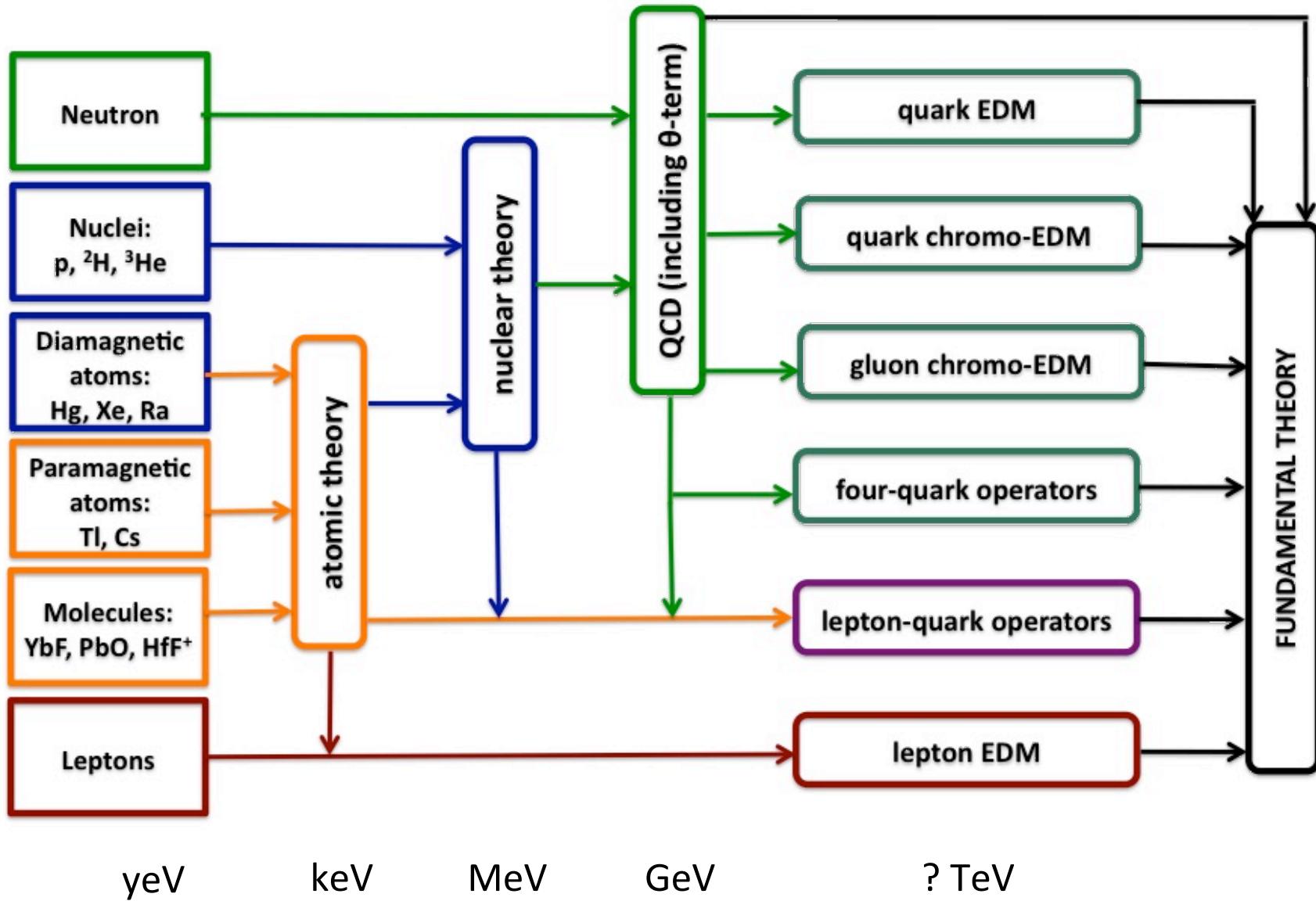
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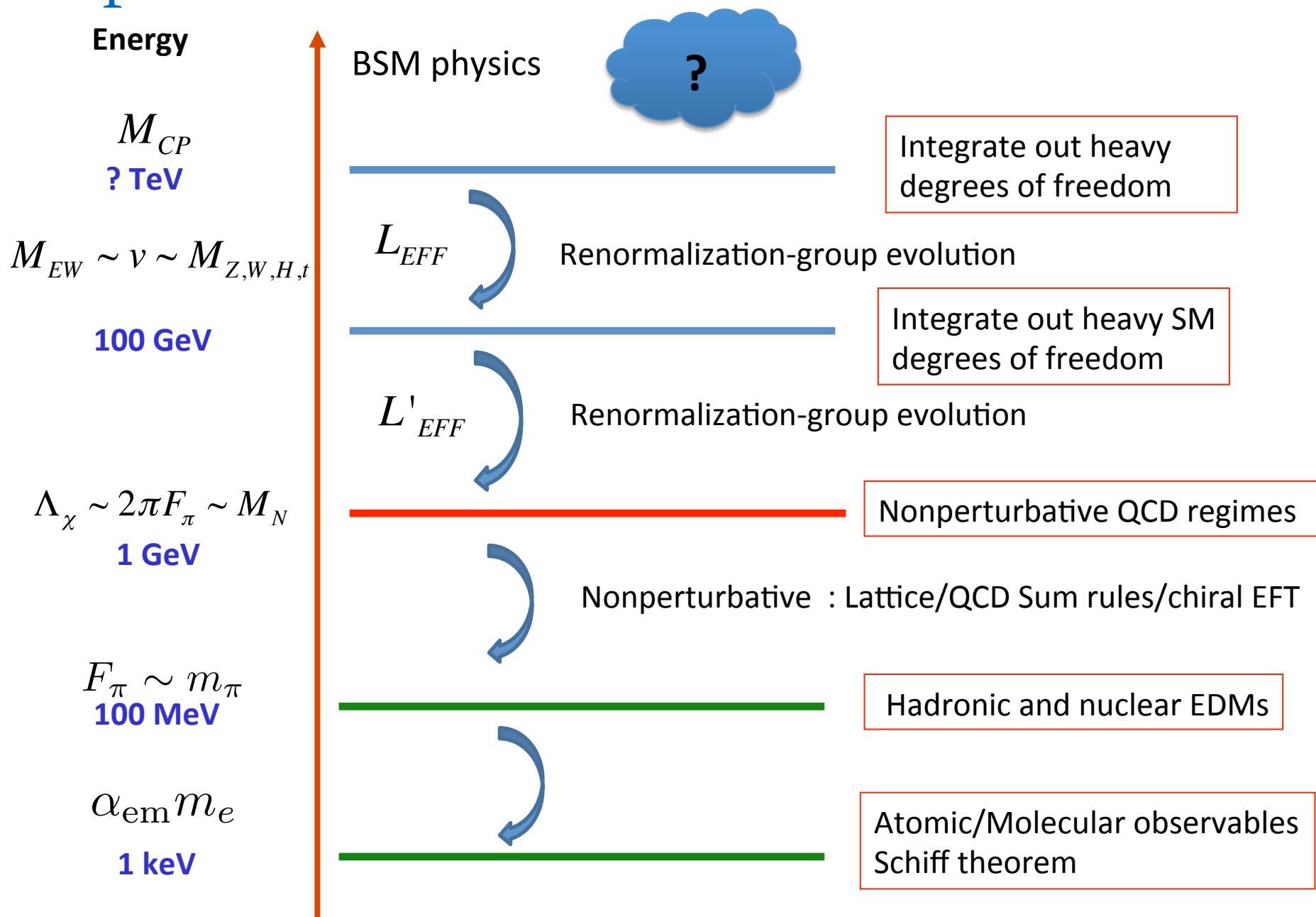
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(Model dependent!)

The EDM landscape



Separation of scales



Step 1: SM as an EFT

- Assume any BSM physics lives at scales $\gg M_{EW}$
- Match to full set of CP-odd operators (model independent *)
 - 1) Degrees of freedom: Full SM field content
 - 2) Symmetries: Lorentz, **SU(3)xSU(2)xU(1)**

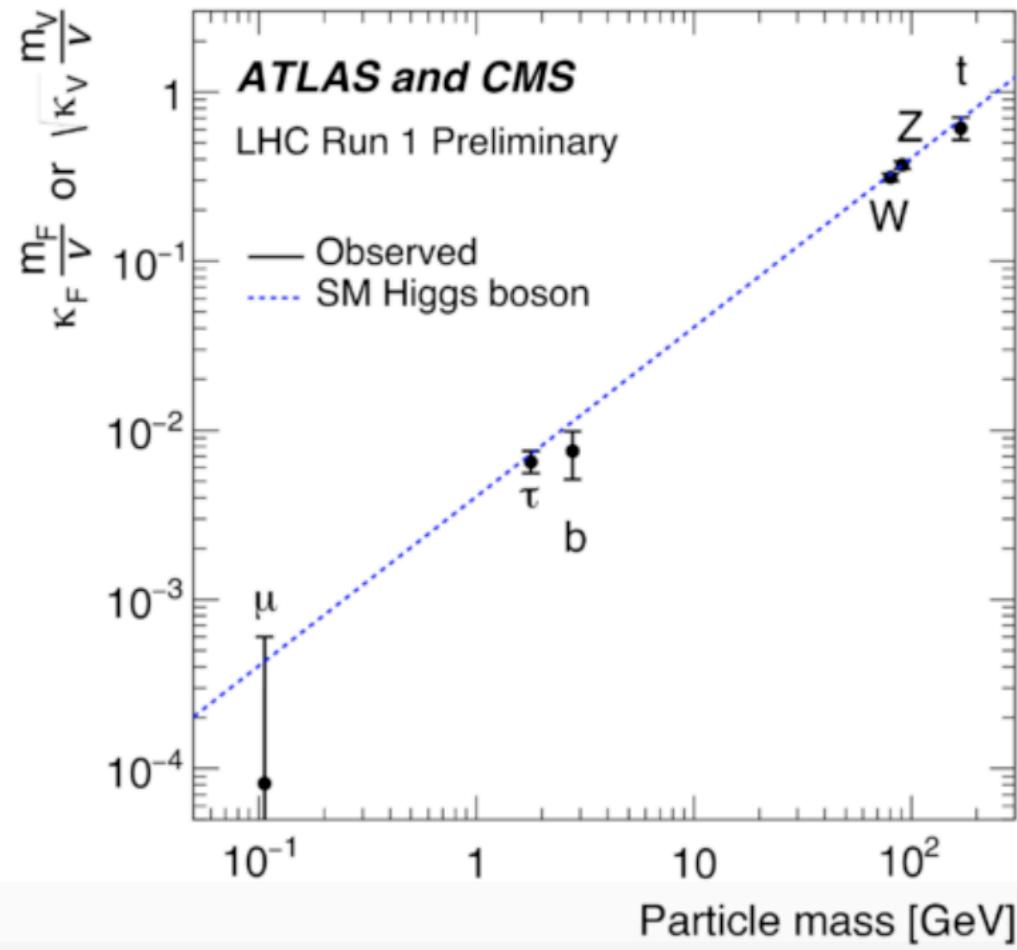
$$L_{new} = \cancel{\frac{1}{M_{CP}} L_5 + \frac{1}{M_{CP}^2} L_6 + \dots}$$

dim-5 generates neutrino masses/mixing, neglected here

* **Big assumption:** no new light fields

One must focus

- Focus on CPV quark-Higgs and gluon-Higgs terms
- Still room for SM deviations
- Typical CP-violation in 2HDMs (and similar models)
- Popular for baryogenesis (e.g. $\text{Im } y_t > 0.1$)
- **Illustrates EDM/LHC complementarity + caveats**



Some earlier studies: Kamenik et al '12, Brod et al '13

Flavor-changing Yukawa's:
Harnik et al '12 & Blankenburg et al '12

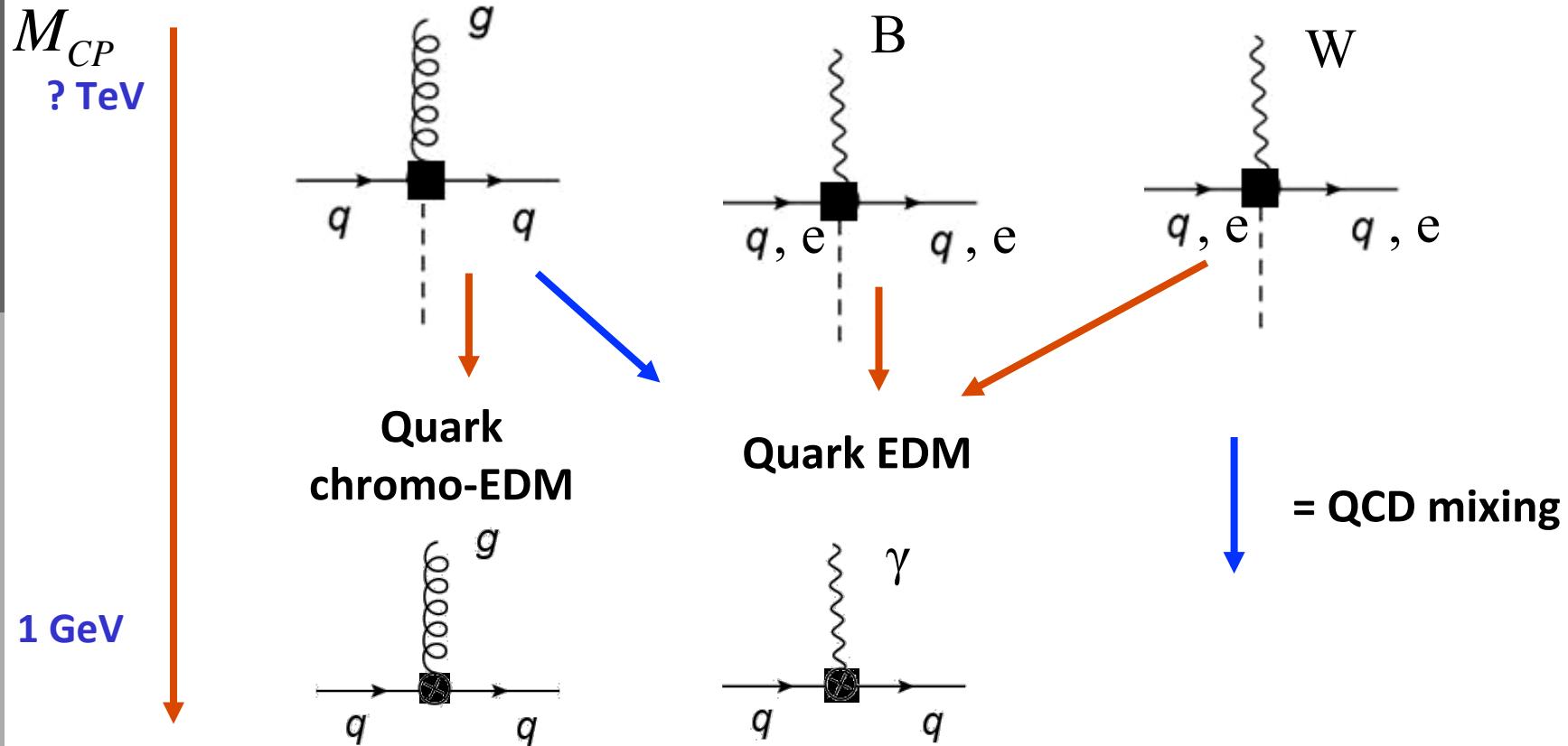
Dipole operators

Requires Higgs: $\Gamma_X \bar{\Psi}_L \sigma^{\mu\nu} \Psi_R X_{\mu\nu} \varphi + h.c.$

X=W,B,G quarks
X=W,B leptons

In most models: $\Gamma_X \propto \frac{m_\Psi}{v M_{CP}^2}$

EDMs typically scale with mass !



Dipole operators

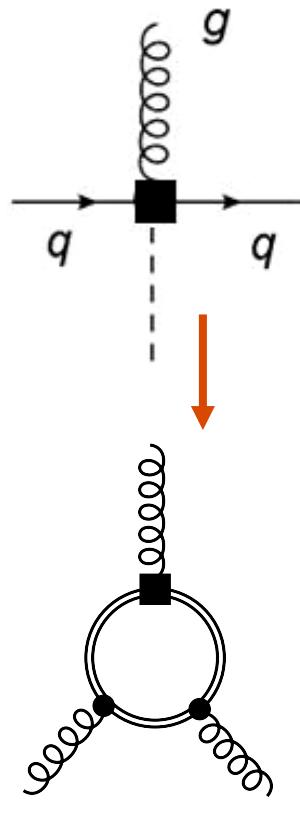
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M_{CP}
? TeV

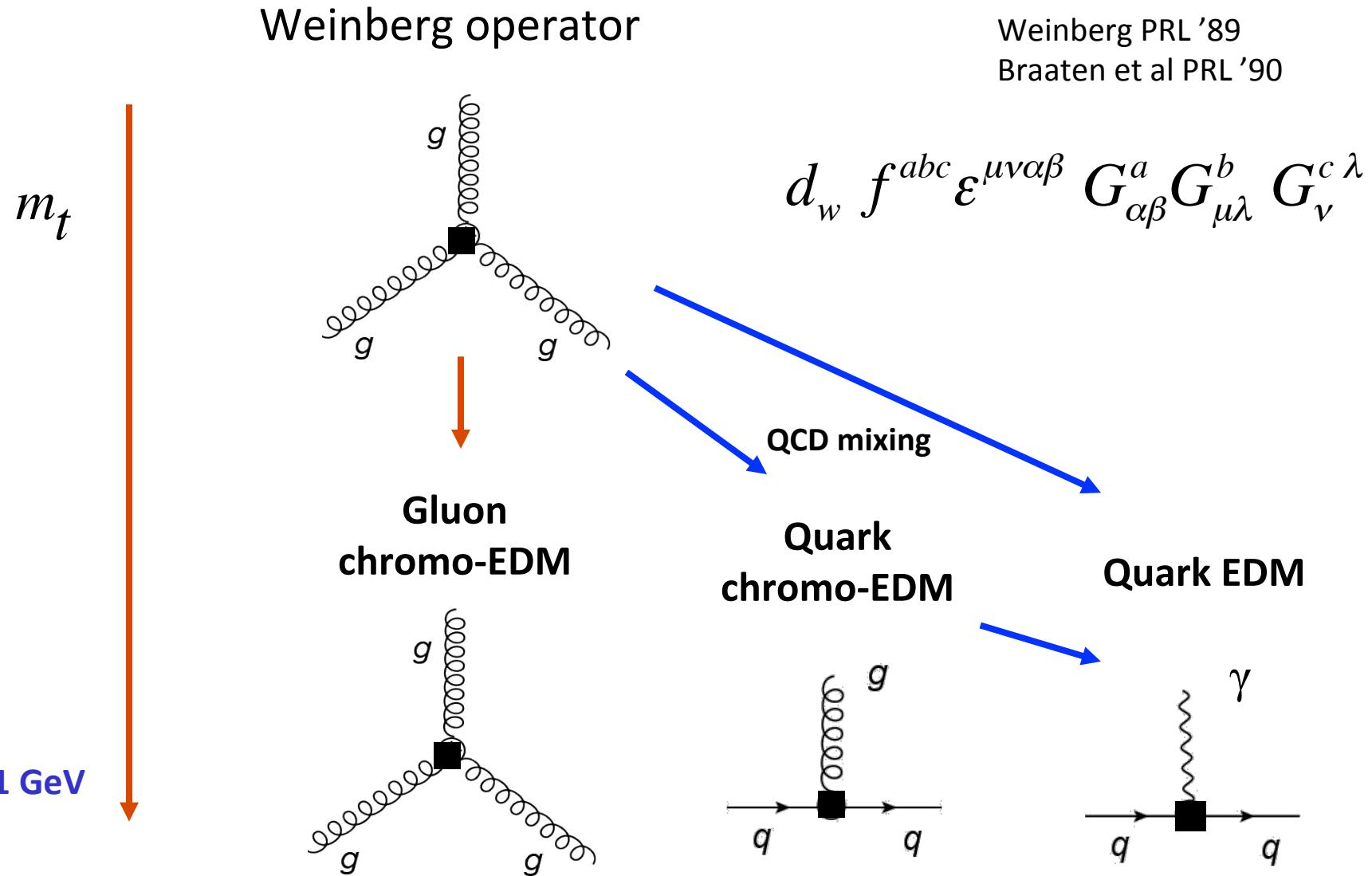


Integrate out heavy quarks

Weinberg operator

$$d_w f^{abc} \epsilon^{\mu\nu\alpha\beta} G_{\alpha\beta}^a G_{\mu\lambda}^b G_{\nu}^{c\lambda}$$

Gluon chromo-EDM



Gluon chromo-EDM

Gauge-Higgs operators (focus on gluonic term)

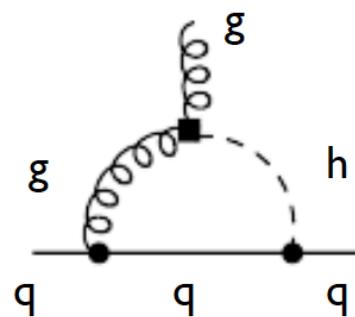
M_{CP}
? TeV

$< M_W$

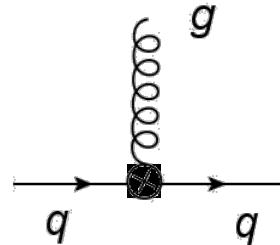
1 GeV

$$\theta' G_{\mu\nu} \tilde{G}^{\mu\nu} (\varphi^\dagger \varphi)$$

Absorb v^2 term in SM (**Strong CP problem**)



Quark
chromo-EDM



Solution of the RG equations

$$\frac{d_q}{m_q}(1 \text{ GeV}) = 1.4 \cdot 10^{-4} Q_q \theta'(1 \text{ TeV})$$

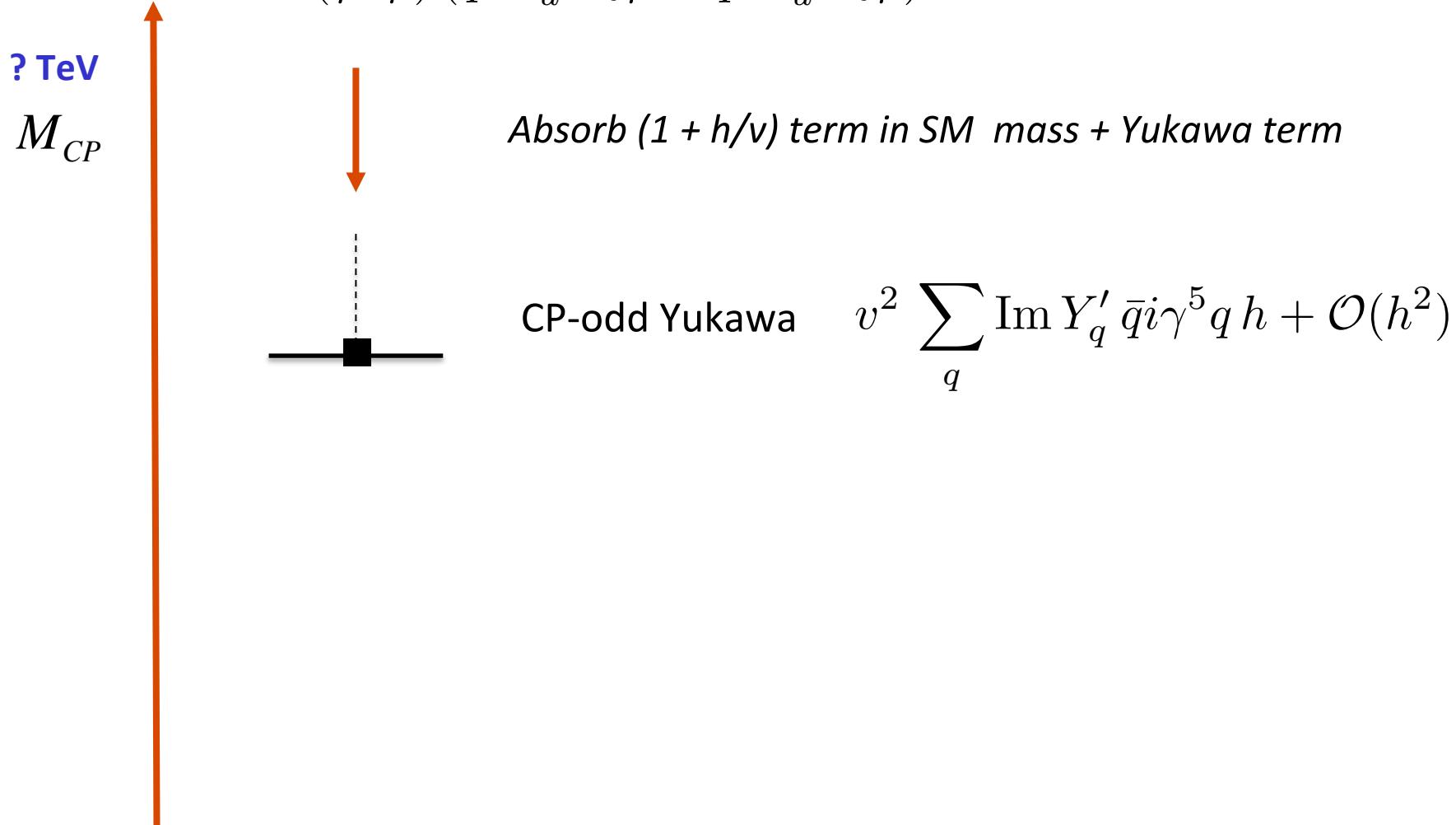
$$\frac{\tilde{d}_q}{m_q}(1 \text{ GeV}) = 1.7 \cdot 10^{-4} \theta'(1 \text{ TeV})$$

$$d_W(1 \text{ GeV}) = -7.3 \cdot 10^{-6} \theta'(1 \text{ TeV})$$

Dim-6 Yukawa interactions

$$\varphi^T = (0 \ h + v)/\sqrt{2}$$

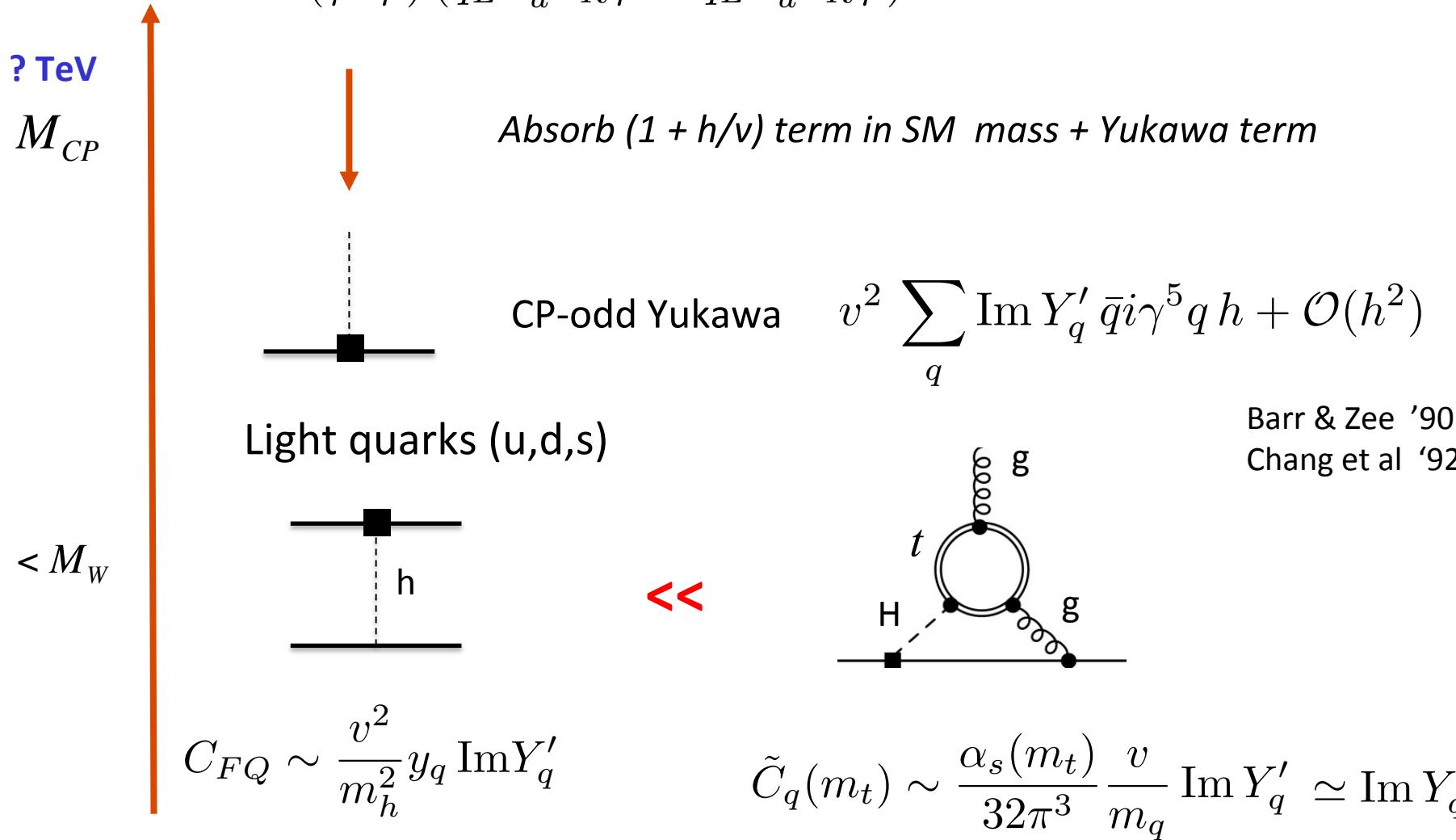
$$-\sqrt{2}(\varphi^\dagger \varphi) (\bar{q}_L Y'_u u_R \tilde{\varphi} + \bar{q}_L Y'_d d_R \varphi) + \text{h.c.}$$



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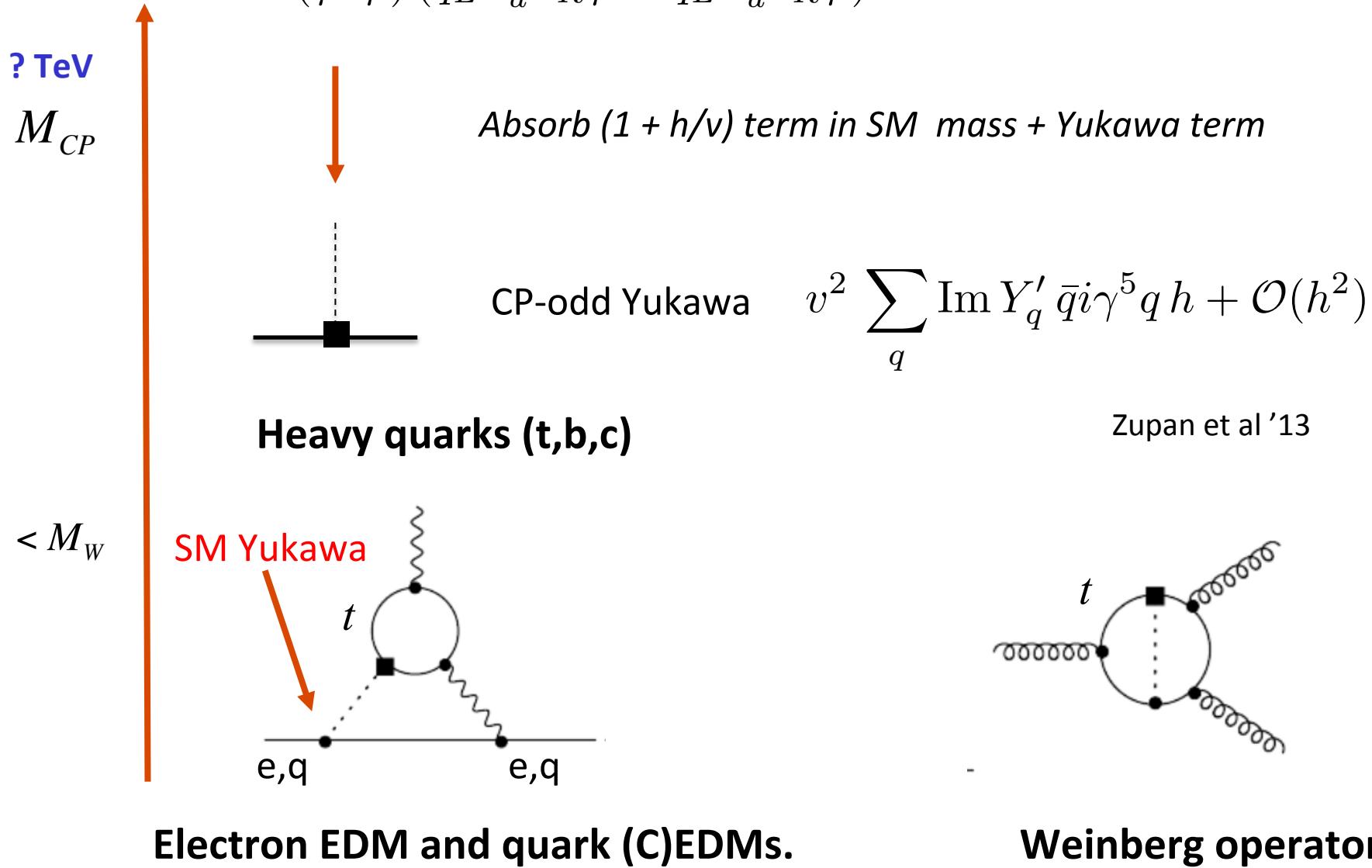
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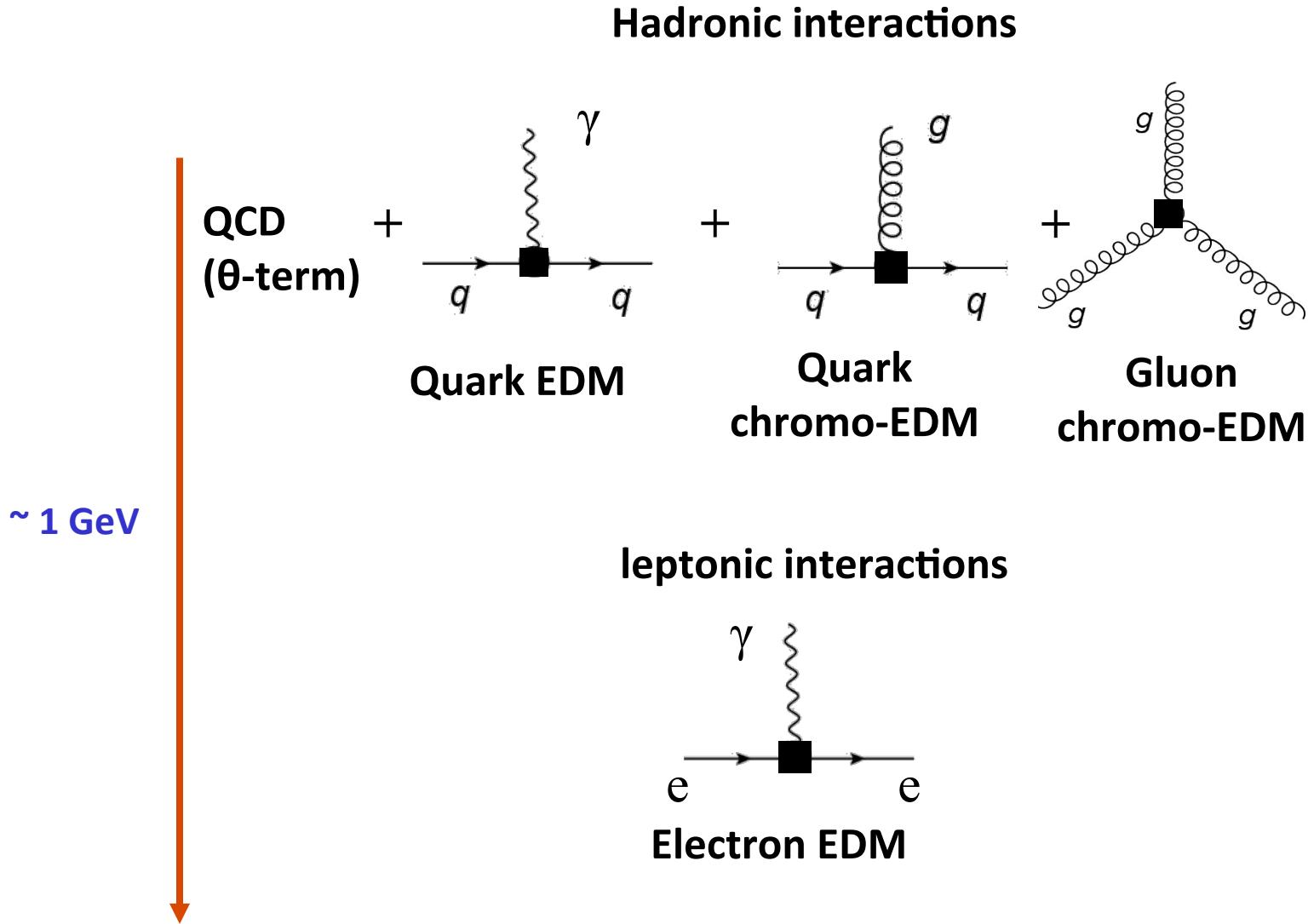
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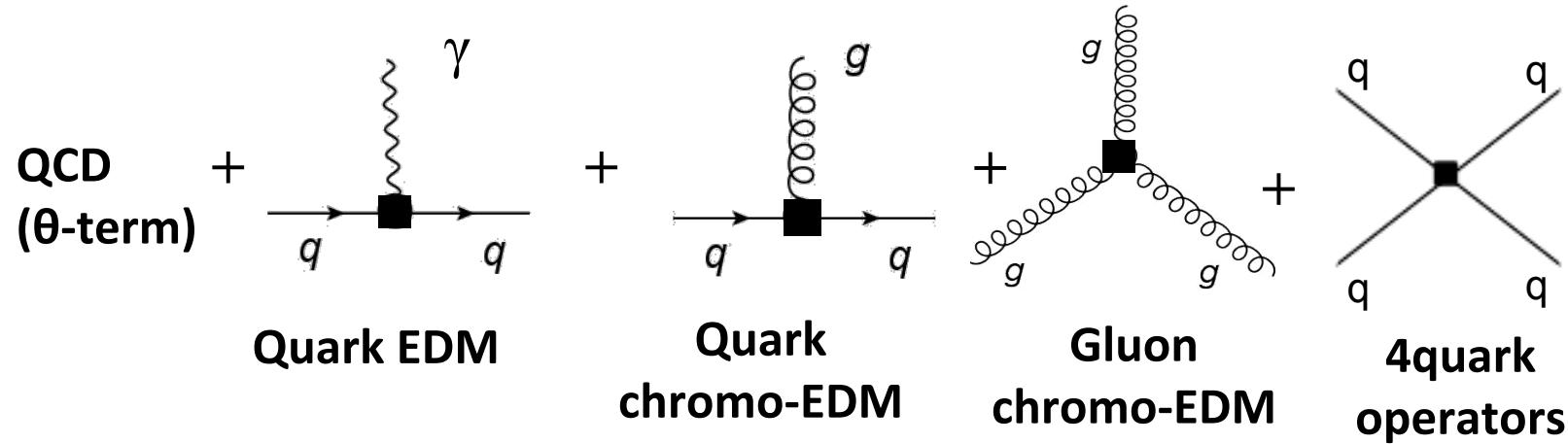


When the dust settles....



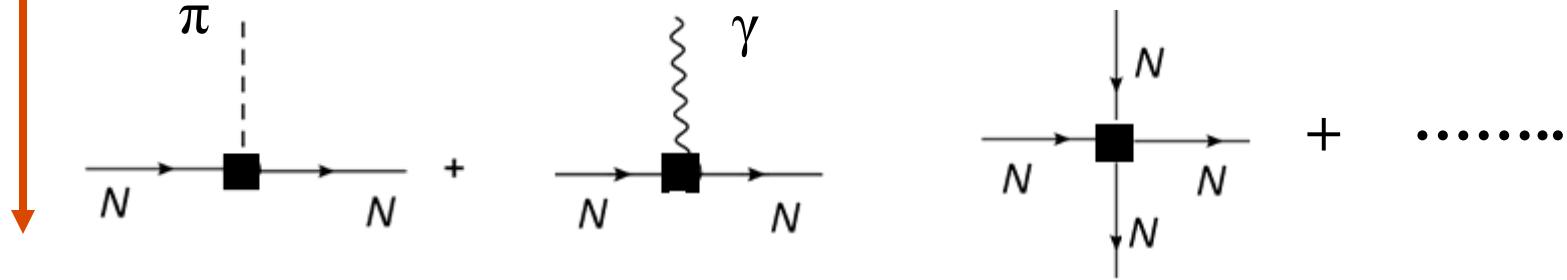
Crossing the barrier

Few GeV



Hadronic/Nuclear/Atomic CP-violation

100 MeV



Chiral EFT

- Use the symmetries of QCD to obtain **chiral Lagrangian**

$$L_{QCD} \rightarrow L_{chiPT} = L_{\pi\pi} + L_{\pi N} + L_{NN} + \dots$$

- Quark masses = 0 \rightarrow QCD has $SU(2)_L \times SU(2)_R$ symmetry
 - Spontaneously broken to $SU(2)$ -isospin
 - Pions are Goldstone bosons
 - Explicit breaking (quark mass) \rightarrow pion mass
- ChPT gives systematic expansion in $Q/\Lambda_\chi \sim m_\pi/\Lambda_\chi \quad \Lambda_\chi \cong 1 \text{ GeV}$
 - **Form of interactions fixed by symmetries**
 - Each interactions comes with an unknown constant (LEC)
 - Successful nucleon-nucleon potential (**chiral EFT**)

The strong CP problem

$$\mathcal{L} = -e^{i\rho} \bar{q}_L M q_R - \theta \frac{\alpha_s}{16\pi} G \tilde{G} + \text{h.c.}$$

After axial U(1) and SU(2) rotations, **two-flavored** mass part of QCD:

$$\mathcal{L} = -\bar{m} \bar{q} q - \varepsilon \bar{m} \bar{q} \tau^3 q + m_\star \bar{\theta} \bar{q} i \gamma^5 q$$

Crewther et al' 79

Baluni '79

$$\bar{\theta} = \theta + 2\rho$$

$$\bar{m} = \frac{m_u + m_d}{2}$$

$$\varepsilon = \frac{m_u - m_d}{m_u + m_d}$$

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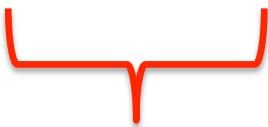
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Linked via $SU_A(2)$ rotation

Crewther et al' 79
Baluni '79

Isospin breaking related to strong CP violation

$$\rho_\theta = -\frac{m_\star \bar{\theta}}{\varepsilon \bar{m}} \simeq -\frac{1 - \varepsilon^2}{2\varepsilon} \bar{\theta}$$

The strong CP problem

$$\mathcal{L} = -\varepsilon \bar{m} \bar{q} \tau^3 q + m_\star \bar{\theta} \bar{q} i \gamma^5 q$$

Explicit ChPT construction shows a relation between:

$$\mathcal{L} = \frac{\delta m_N}{2} \bar{N} \tau^3 N + \bar{g}_0 \bar{N} \pi \cdot \tau N \quad N = (p \ n)$$

Nucleon mass splitting
(strong part, no QED part!)



CP-odd pion-nucleon interaction

The strong CP problem

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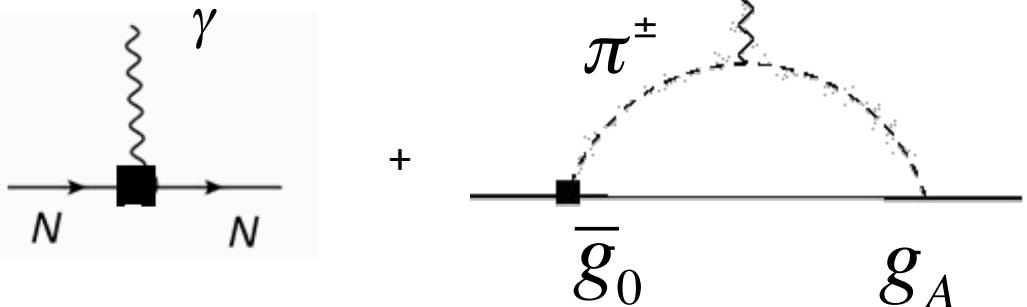
**CP-odd pion-nucleon
interaction**

$$\frac{\bar{g}_0}{f_\pi} = \delta m_N \rho_\theta = -\delta m_N \frac{1 - \varepsilon^2}{2\varepsilon} \bar{\theta} = -(15.5 \pm 2.5) \cdot 10^{-3} \bar{\theta}$$

- Using **lattice results** for (nucleon, quark) mass differences
Walker-Loud '14, Borsanyi '14, Aoki (FLAG) '13,
- This and other relations hold up to N2LO in SU(2) and SU(3) ChPT
JdV, Mereghetti, Walker-Loud '15

The strong CP problem

Nucleon EDM



$$d_n = \bar{d}_0 - \bar{d}_1 - \frac{e g_A \bar{g}_0}{4\pi^2 F_\pi} \left(\ln \frac{m_\pi^2}{M_N^2} - \frac{\pi}{2} \frac{m_\pi}{M_N} \right)$$

CP-even pion-nucleon coupling

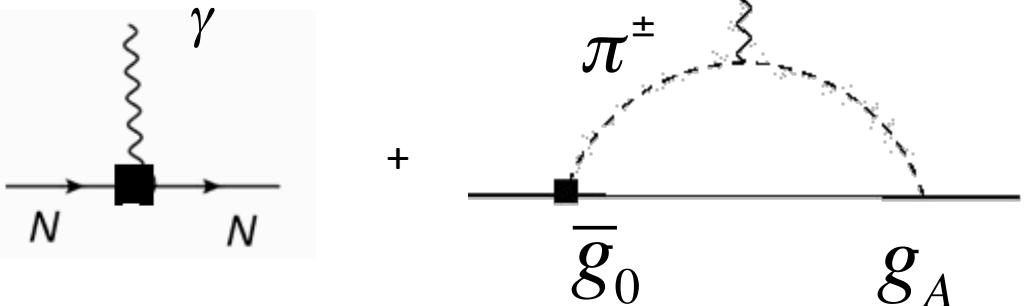
- Loop is divergent.... Need a counter-term
- Loop is **enhanced** by chiral logarithm (long-rang physics)

$$\bar{g}_0 = -(15.5 \pm 2.5) \cdot 10^{-3} \bar{\theta} \quad \longrightarrow \quad d_n \simeq -2.5 \cdot 10^{-16} \bar{\theta} \text{ e cm}$$

• Experimental constraint: $\longrightarrow \bar{\theta} < 10^{-10}$

The strong CP problem

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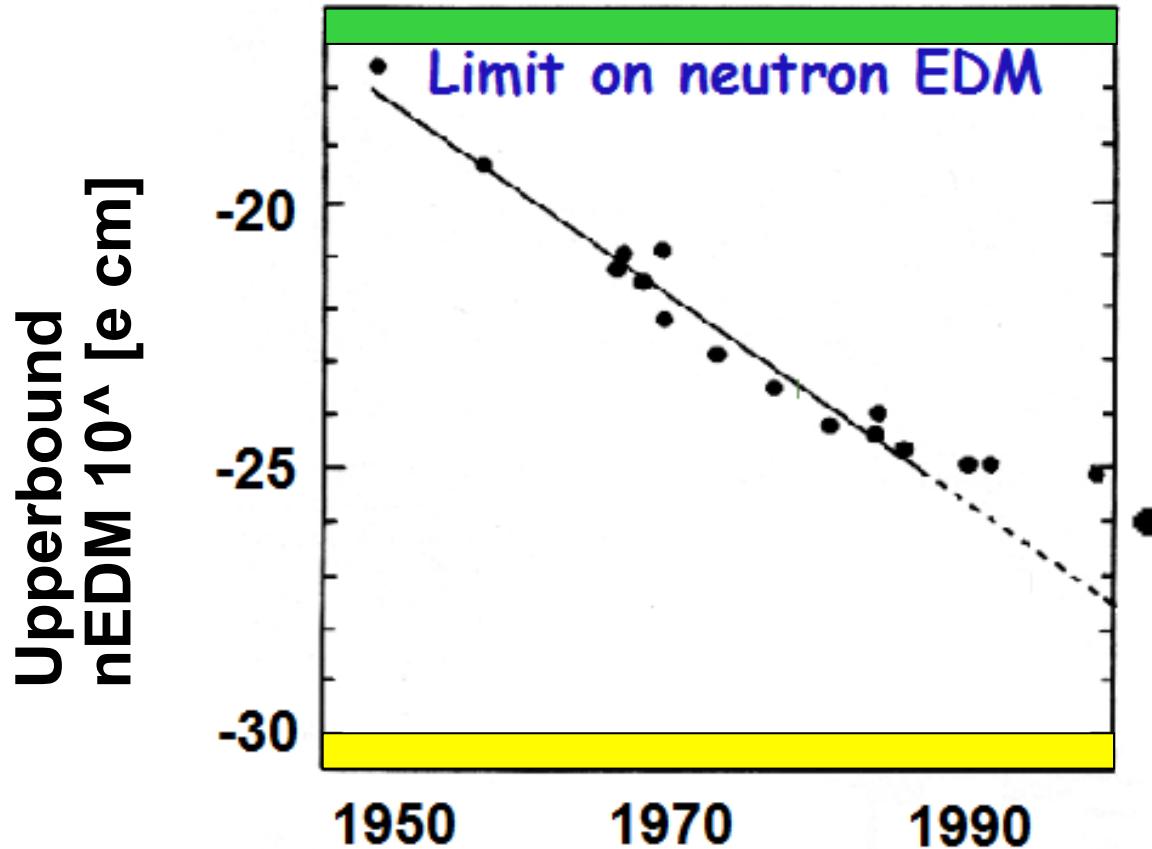
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- Lattice + ChPT $d_n = -(2.7 \pm 1.2) \cdot 10^{-16} \bar{\theta} e$ Shintani et al '12

- $d_n = -(3.9 \pm 1.0) \cdot 10^{-16} \bar{\theta} e$ Guo et al '15

Strong CP problem



- Set one quark mass to zero. Disfavored by lattice QCD.
- Assume P or CP exact at high energies (e.g. left-right models)
- **Peccei-Quinn mechanism \rightarrow axions (we assumed this)**

Left with the dim6 sources

- ❖ Quark EDM accurately determined recently !

T. Bhattacharya et al '15

$$d_n = -(0.22 \pm 0.03)d_u + (0.74 \pm 0.07)d_d + (0.008 \pm 0.01)d_s$$

- ChPT extrapolation to **physical pion mass**

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- ChPT extrapolation to **physical pion mass**

- ❖ Quark CEDM no lattice calculations yet. **But in progress.**

ChiPT/QCD sum rules: pion-nucleon couplings and nucleon EDMs

50-75% uncertainty

Pospelov, Ritz '02 '05

JdV et al '10 '13

Hisano et al '12 '13

- ❖ Weinberg **estimate for nEDM**

$$d_n = \pm [(50 \pm 40) \text{ MeV}] e d_W$$

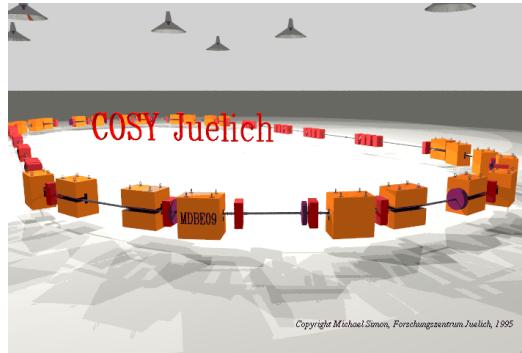
Demir et al '03

pion-nucleon couplings **suppressed** (chiral symmetry)

Storage rings experiments

Farley *et al* PRL '04

- New kid on the block: **Charged particle in storage ring**

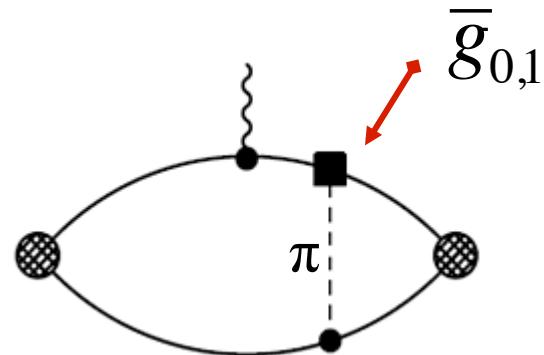
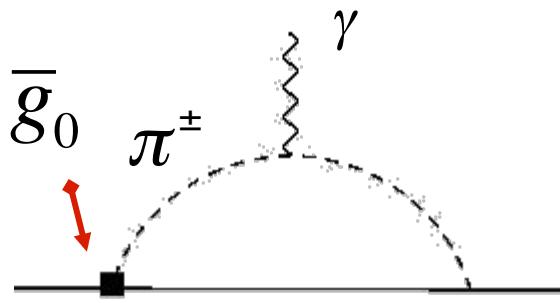


Bennett *et al* (BNL g-2) PRL '09

- Limit on muon EDM $d_\mu \leq 1.8 \cdot 10^{-19} \text{ e cm}$ (95% C.L.)
- Proposals to measure EDMs of proton, deuteron, ^3He at level $\sim 10^{-29} \text{ e cm}$
- Other light nuclei?

COSY @ Jülich
 Brookhaven/Fermilab

Why light nuclei?



- **Tree-level:** no loop suppression
- Very good **theoretical control** !

$$d_A = \langle \Psi_A | \vec{J}_{CP} | \Psi_A \rangle + 2 \langle \Psi_A | \vec{J}_{CP} | \tilde{\Psi}_A \rangle$$

$$(E - H_{PT}) |\Psi_A\rangle = 0 \quad (E - H_{PT}) |\tilde{\Psi}_A\rangle = V_{CP} |\Psi_A\rangle$$

Input

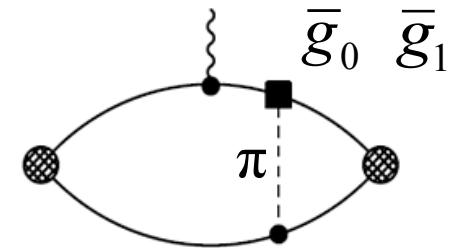
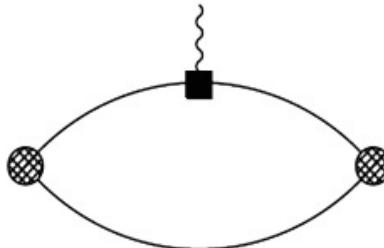
1. CP-even and -odd potential from **chiral EFT**
2. Solve Schrodinger equations numerically

Epelbaum et al '05
Maekawa et al '11

Example: deuteron EDM

Target of storage ring measurement

- Two contributions (NLO)
 1. Sum of nucleon EDMs
 2. CP-odd pion exchange



Errors from Bsaisou et al JHEP '14

$$d_D = d_n + d_p + [(0.18 \pm 0.02) \bar{g}_1 + (0.0028 \pm 0.0003) \bar{g}_0] e \text{ fm}$$

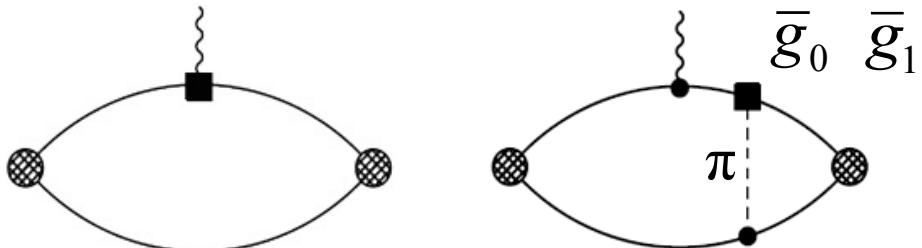
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(chiral corrections + cut-off dependence)

Strong isospin filter

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Strong isospin filter

- Tree-level pion exchange can dominate the nuclear EDM
- $d_D \sim |6 dn|$ for qCEDMs, $d_D \sim dn + dp$ for qEDM/Weinberg
- Differentiate between various BSM models (2HDM, MLRSM)

Diamagnetic EDMs

Strongest bound on atomic EDM:

$$d_{^{199}Hg} < 3.1 \cdot 10^{-29} e\text{ cm}$$

New measurements expected: Hg, Ra , Xe,

Schiff Theorem: EDM of nucleus is screened by electron cloud if:

1. Point particles
2. Non-relativistic kinematics

Schiff, '63

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Screening incomplete: nuclear finite size (Schiff moment \mathbf{S})

Typical suppression:
$$\frac{d_{Atom}}{d_{nucleus}} \propto 10 Z^2 \left(\frac{R_N}{R_A} \right)^2 \approx 10^{-3}$$

- **Atomic** part well under control

$$d_{^{199}Hg} = (2.8 \pm 0.6) \cdot 10^{-4} S_{Hg} e \text{ fm}^2$$

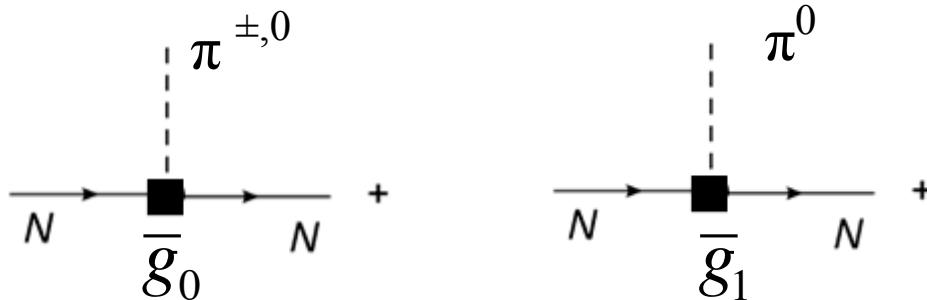
Dzuba et al, '02, '09

$$d_{^{225}Ra} = (7.2 \pm 1.5) \cdot 10^{-4} S_{Ra} e \text{ fm}^2$$

Sing et al, '15

Calculating Schiff Moments

Task: Calculate Schiff Moments of Hg, Ra, Xe, ...



Flambaum, de Jesus,
Engel, Dobaczewski,
Dmitriev, Sen'kov,.....

- **Typically only one-pion exchange** (sometimes nucleon EDMs)
- **Very complicated** many-body calculation Dmitriev, Sen'kov '03
- Use nuclear model and mean-field theory

$$S_{\text{Hg}} = [(0.35 \pm 0.3)\bar{g}_0 + (0.35 \pm 0.70)\bar{g}_1] e \text{ fm}^3$$

- **Large** uncertainties. Even unknown sign !

Probing the electron EDM

Bound on Tl EDM

$$d_{^{205}Tl} < 9 \cdot 10^{-25} \text{ e cm}$$

Regan et al '02

What about screening? Schiff theorem violated by **relativity**

$$d_A(d_e) = K_A d_e \quad K_A \propto Z^3 \alpha_{em}^2$$

Sandars '65

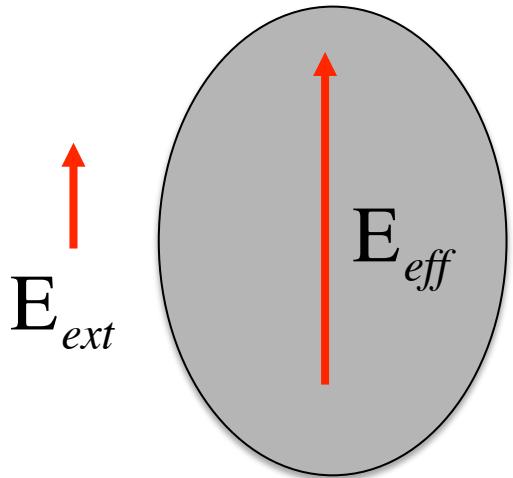
Strong enhancement!

$$K_{Tl} = -(570 \pm 20) \longrightarrow d_e < 1.6 \cdot 10^{-27} \text{ e cm}$$

Polar molecules

Polar molecules: Convert small external to huge internal field

Kozlov et al '94 '97, Quiney et al '98, Mayer, Bohn '08



$$\Delta E_{YbF} = (15 \pm 2) \cdot GeV \left(\frac{d_e}{e\ cm} \right)$$

$$\Delta E_{ThO} = (80 \pm 10) \cdot GeV \left(\frac{d_e}{e\ cm} \right)$$

Meyer, Bohn '08, Skipnikov et al '13, Fleig, Nayak '14,

The most recent constraint is then:

$$d_e < 8.7 \cdot 10^{-29} \ e\ cm$$

One order improvement expected next 5 years...

Baron et al '13

Strategy for setting limits

Study impact of uncertainties in the hadronic/nuclear EDMs

1. **Central:** use central value matrix elements (most common method)
2. **RFit** (“Range-Fit”): vary matrix elements in their allowed ranges; minimized chi-squared (=most conservative bounds)

Strategy copied from CKMfitter group '04

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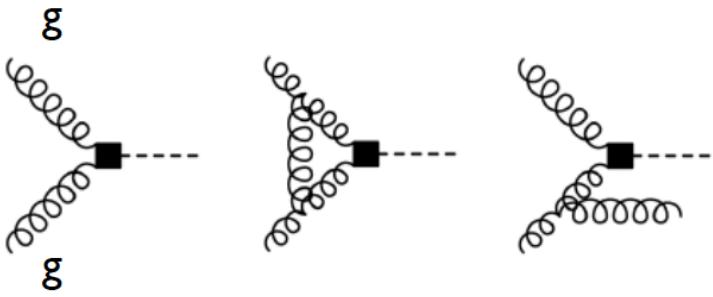
3. **RFit+:** Rfit with improved theory (50% uncertainties everywhere)

Realistic (but challenging) target for Lattice-QCD + nuclear structure

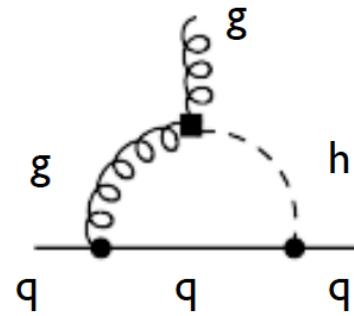
- **Dedicated Amherst workshop, January '15 → road map**
“Hadronic Matrix Elements for Probes for CP-violation”

Anomalous gluon-higgs coupling

LHC: Higgs production via gluon fusion



Low Energy: quark (C)EDM + Weinberg



$$\begin{aligned}\mu &= \frac{\sigma_{\text{GGF}}^{SM} + \sigma_{\text{GGF}}^{\theta'}}{\sigma_{\text{GGF}}^{SM}} \\ &= 1 + (2.28 \pm 0.1)(v^2 \theta')^2\end{aligned}$$

- Cross section known to N2LO
- Error from scale variation + PDFs

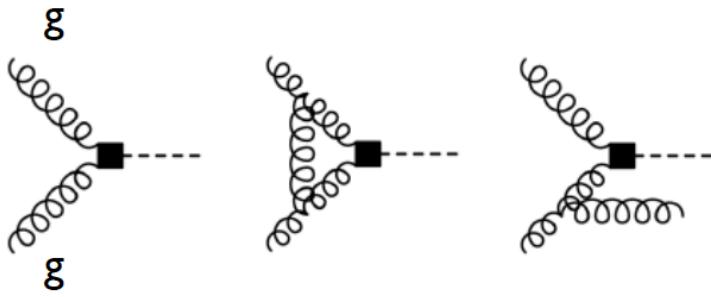
Harlander, Kilgore '02, '03
Anastasiou, Melnikov '02 '03

$$\begin{aligned}\frac{d_q}{m_q}(1 \text{ GeV}) &= 1.4 \cdot 10^{-4} Q_q \theta'(1 \text{ TeV}) \\ \frac{\tilde{d}_q}{m_q}(1 \text{ GeV}) &= 1.7 \cdot 10^{-4} \theta'(1 \text{ TeV})\end{aligned}$$

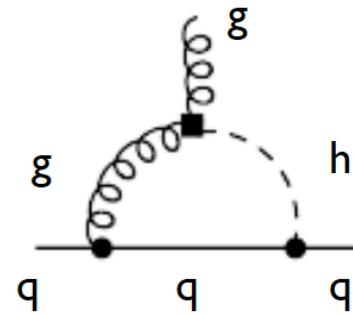
$$d_W(1 \text{ GeV}) = -7.3 \cdot 10^{-6} \theta'(1 \text{ TeV})$$

Anomalous gluon-higgs coupling

LHC: Higgs production via gluon fusion



Low Energy: quark (C)EDM + Weinberg



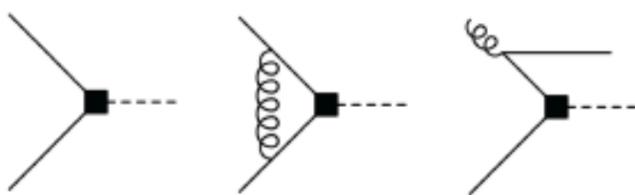
Current experiments

	$v^2\theta'$	d_n	d_{Hg}	d_n, d_{Hg} (comb)	LHC (CMS)
Central	0.06	0.04		0.04	0.27
RFit	0.23	X		0.23	0.27
RFit+				0.05	0.27

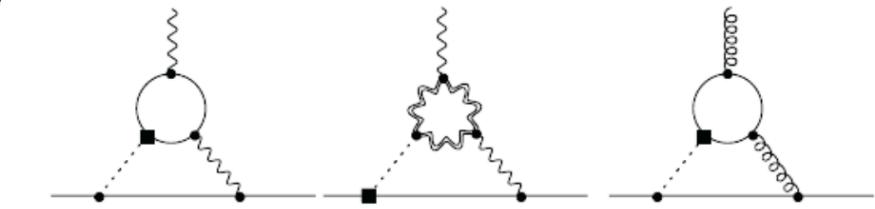
Bounds on couplings at the scale $\mu = M_{BSM} = 1 \text{ TeV}$

Yukawa's u, d, s, c, b

LHC: Higgs production



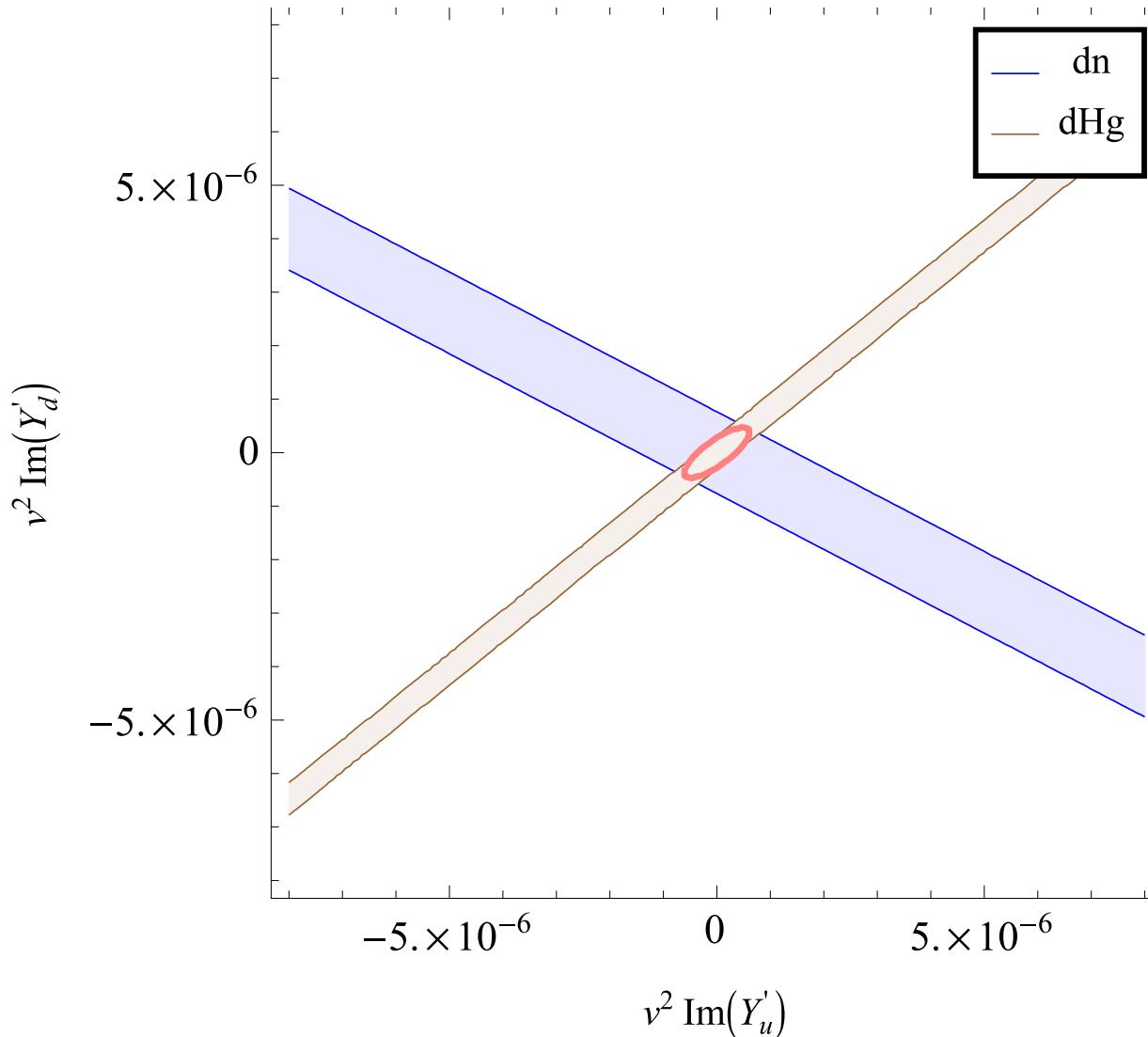
Low Energy: quark (C)EDM, Weinberg, and de



	$v^2 \text{Im } Y'_u$	$v^2 \text{Im } Y'_d$	$v^2 \text{Im } Y'_s$	$v^2 \text{Im } Y'_c$	$v^2 \text{Im } Y'_b$
Central	$3.9 \cdot 10^{-7}$	$3.0 \cdot 10^{-7}$	$4.3 \cdot 10^{-4}$	$1.1 \cdot 10^{-3}$	$8.4 \cdot 10^{-3}$
Rfit	$2.8 \cdot 10^{-6}$	$1.5 \cdot 10^{-6}$	0.42	$6.5 \cdot 10^{-3}$	0.041
LHC	$6.0 \cdot 10^{-3}$	$7.0 \cdot 10^{-3}$	0.020	0.016	0.036

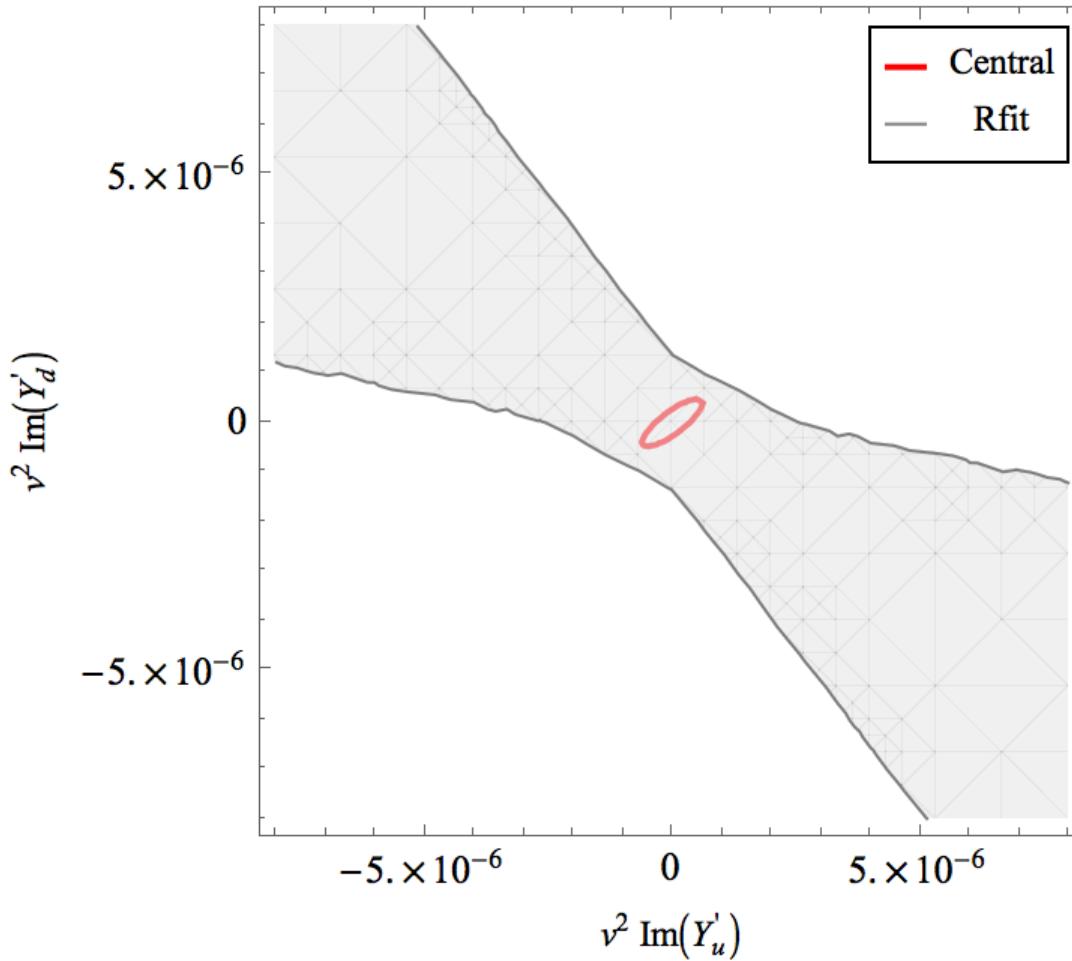
- u- and d-quark **out of reach**, but for s, c, b LHC is **better** or **comparable**
- With improved theory EDMs could beat LHC Higgs production
- **CAVEAT: Only bounds from Higgs production.**
- Should study: Higgs decays and CP-odd correlations

Constraining CPV Yukawa's



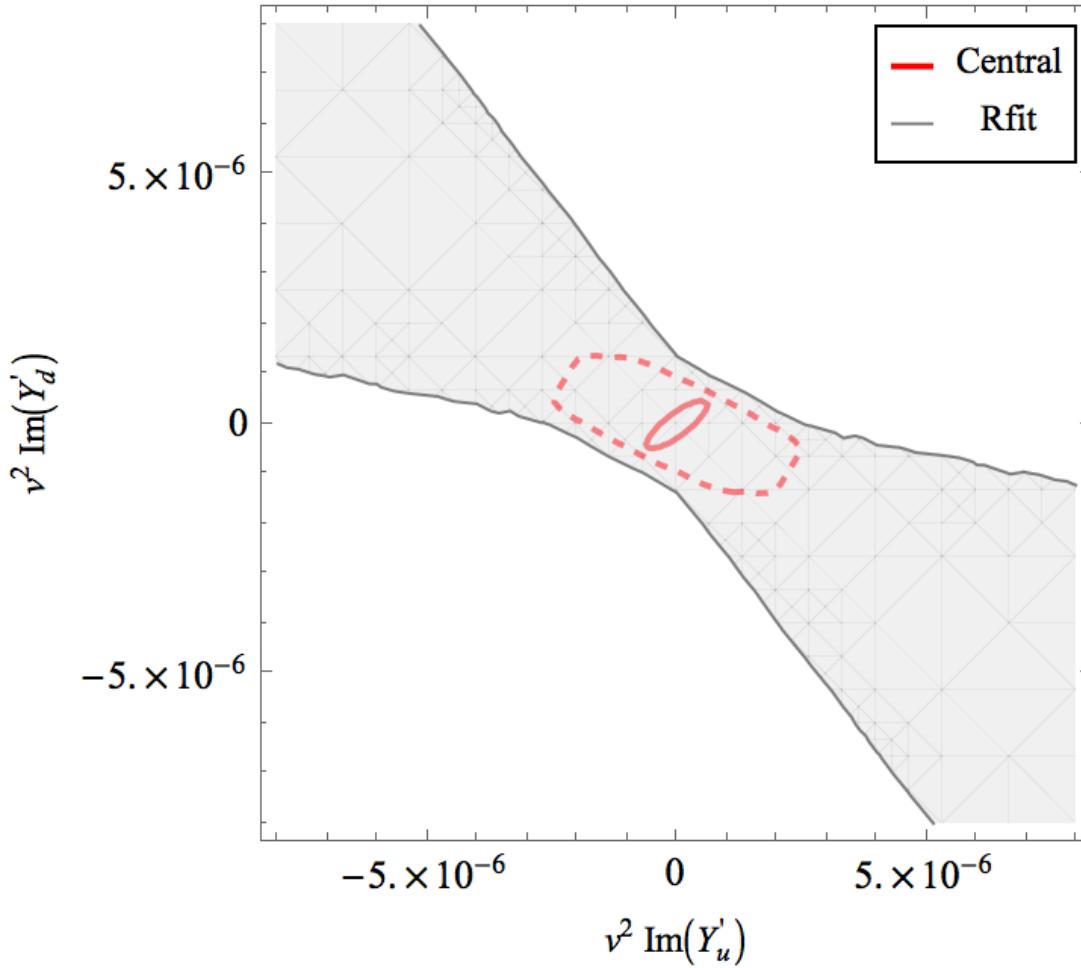
EDMs bound imaginary Yukawa's at **ppm** level

Constraining CPV Yukawa's



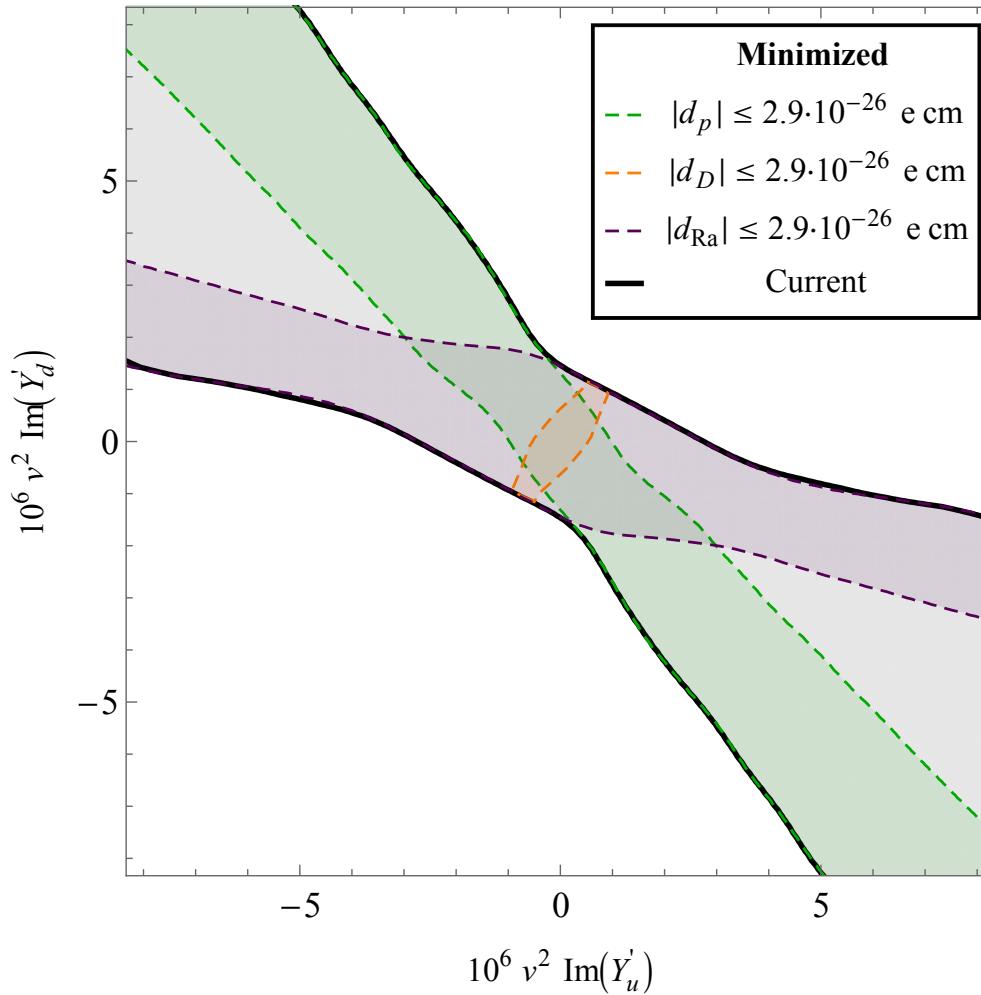
Due to nuclear/hadronic uncertainties a **free direction** emerges

Improved matrix elements



50% matrix elements: almost maximum reach

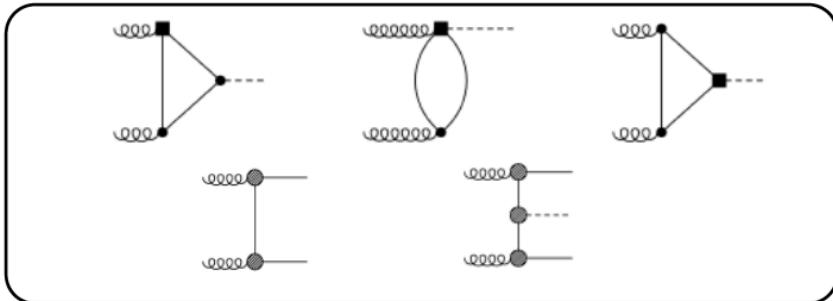
Additional probes



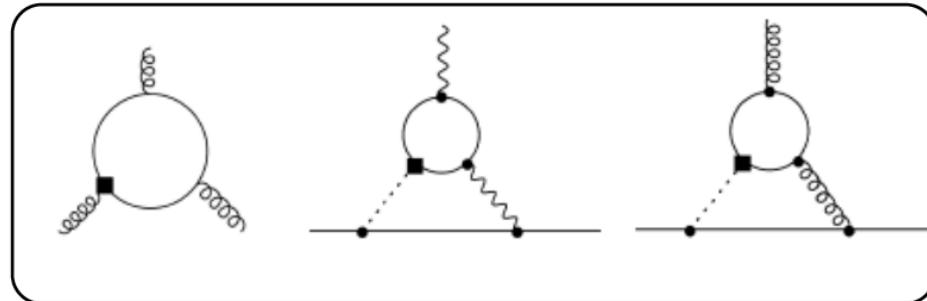
Deuteron EDM **very** complementary !
Radium as well, but uncertainties are larger

Top quark Yukawa and CEDM

LHC: $\text{pp} \rightarrow \text{h}$ (via ggF), $\text{t}\bar{\text{t}}$, tth



Low Energy: quark (C)EDM, Weinberg, and d_e



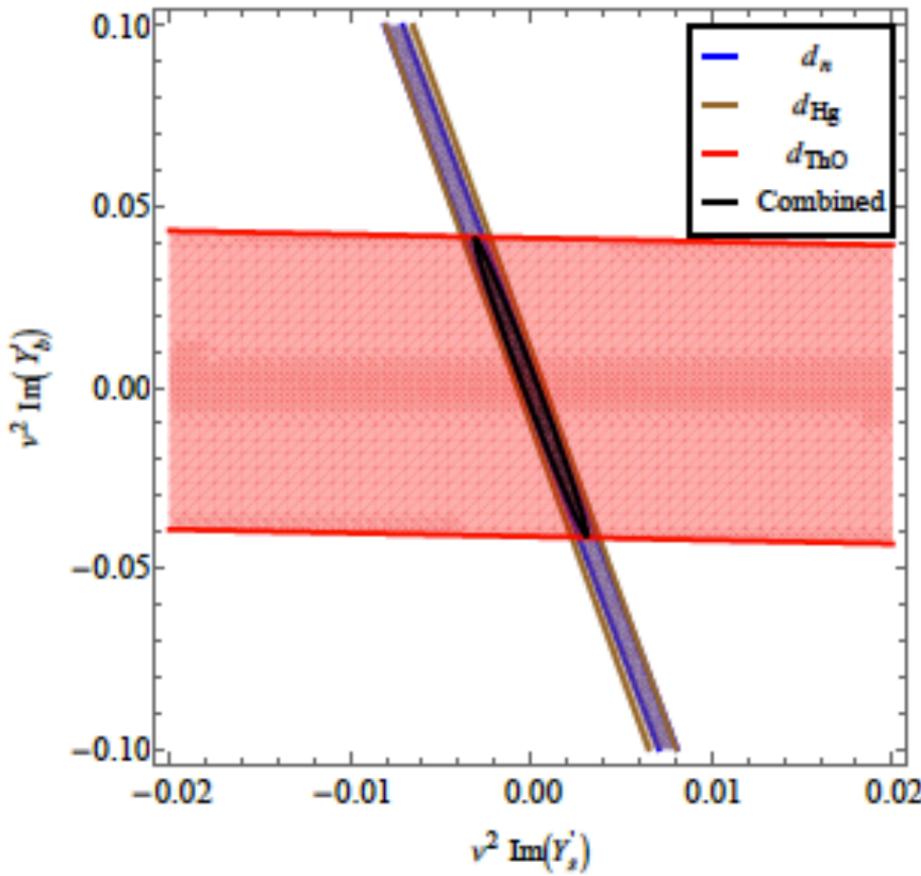
$v^2 \text{Im } Y'_t$	d_n	d_{Hg}	d_e	$[d_n, d_{Hg}, d_e](\text{comb})$	LHC
Central	0.047	0.036	$7.8 \cdot 10^{-3}$	$7.8 \cdot 10^{-3}$	0.15
Rfit	0.11	> 1	$7.8 \cdot 10^{-3}$	$7.8 \cdot 10^{-3}$	0.15

- Richer collider phenomenology (ggFusion, ttbar, ttbar h)
- Strong constraint on $\text{Im } Y_t$ from eEDM ($\sim \text{SM electron Yukawa!}$)
- Little uncertainty on EDM constraint

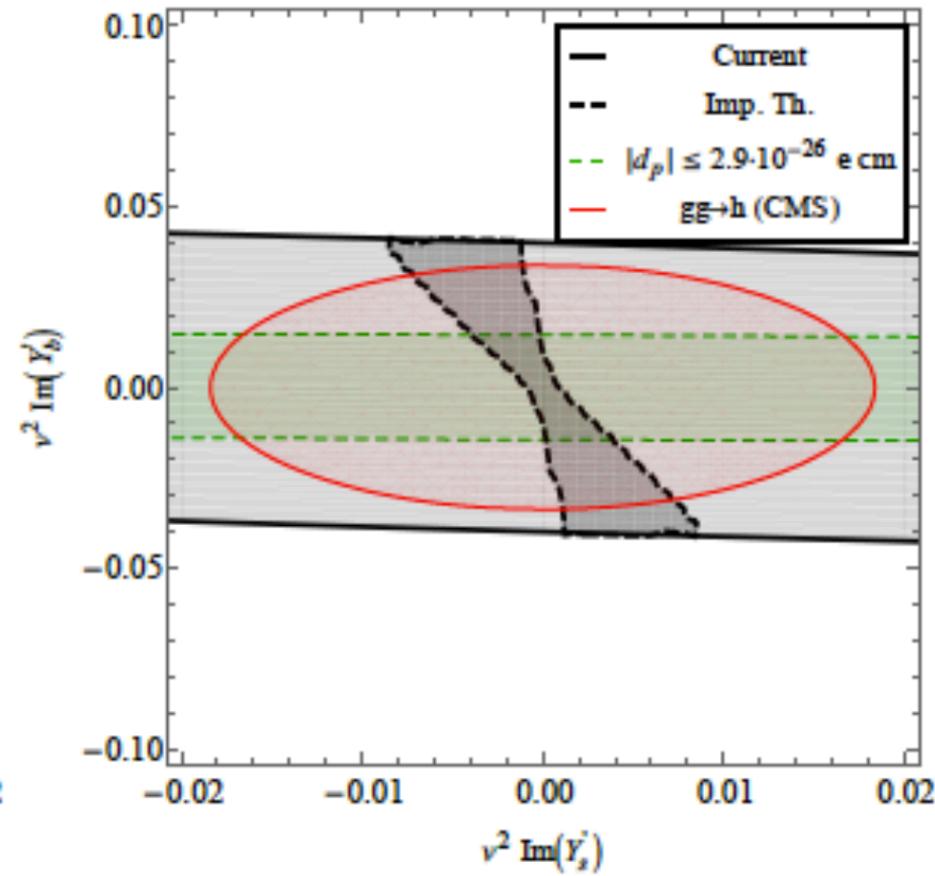
LHC-EDM complementarity

- Two coupling analysis (bottom and strange Yukawa)

Central



RFit

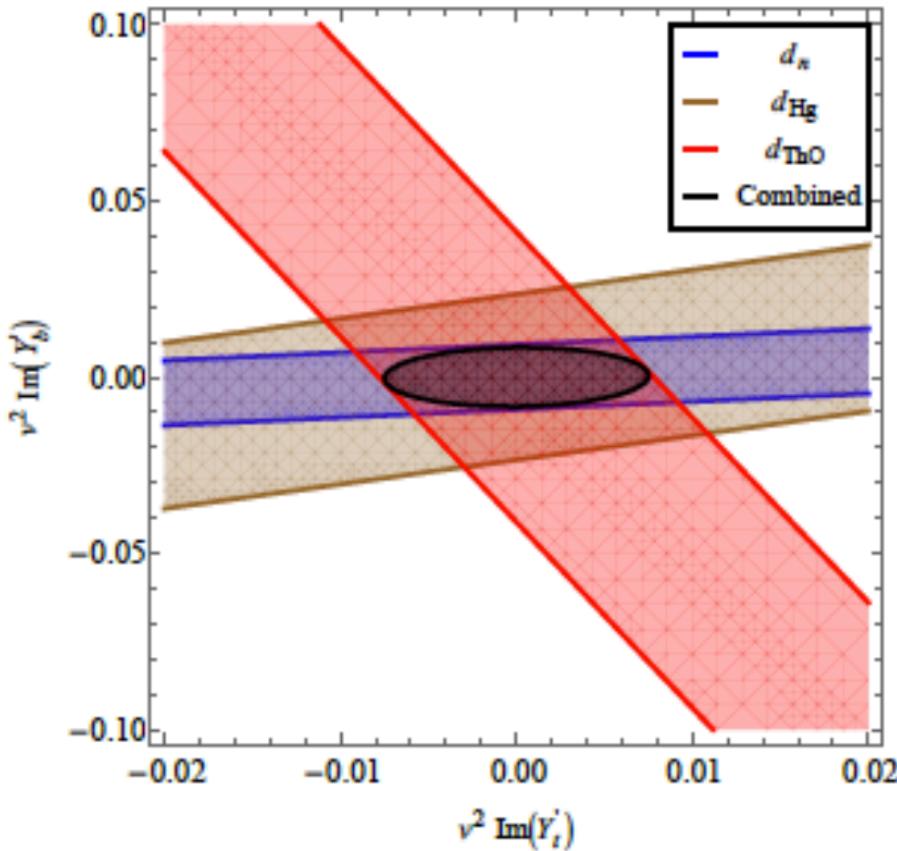


LHC or improved theory removes free direction

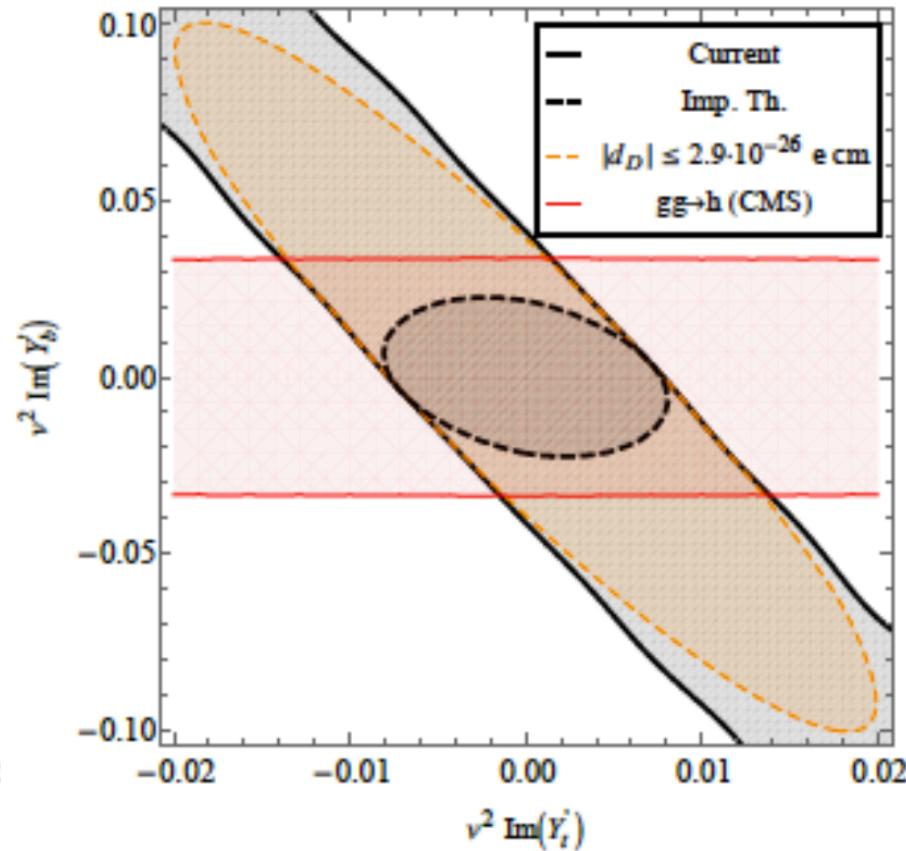
LHC-EDM complementarity

- Two coupling analysis (bottom and top Yukawa)

Central



RFit



LHC or improved theory removes free direction

Summary table

	$v^2 \text{Im } Y'_u$	$v^2 \text{Im } Y'_d$	$v^2 \text{Im } Y'_c$	$v^2 \text{Im } Y'_s$	$v^2 \text{Im } Y'_t$	$v^2 \text{Im } Y'_b$	$v^2 \theta'$
EDMs	$2.8 \cdot 10^{-6}$	$1.5 \cdot 10^{-6}$	$6.3 \cdot 10^{-3}$	0.42	$7.8 \cdot 10^{-3}$	0.041	0.23
LHC Run 1	0.06	0.07	0.02	0.015	0.15	0.038	0.27

- Pseudoscalar Yukawa's in units of SM Yukawa's

$$\mathcal{L} = \frac{m_q}{v} \tilde{\kappa}_q \bar{q} i\gamma_5 q h$$

$\tilde{\kappa}_u$	$\tilde{\kappa}_d$	$\tilde{\kappa}_s$	$\tilde{\kappa}_c$	$\tilde{\kappa}_b$	$\tilde{\kappa}_t$
0.45	0.11	58	2.3	3.6	0.01

- Impressive constraints on up, down, and top !
- Can improve a lot with theory + experimental improvements

Conclusion/Summary

EFT approach

- ✓ Framework exists for CP-violation (EDMs) from 1st principles
- ✓ Keep track of **symmetries** (e.g. chiral) from multi-Tev to atomic scales
- ✓ Specific models can be matched to EFT framework (not discussed here)

CP-violating quark- and gluon-Higgs interactions

- ✓ EDMs and LHC Higgs production are complementary
- ✓ EDMs have a **potential** edge but suffer from hadronic/nuclear uncertainties
- ✓ Set a **target** for lattice/nuclear structure to improve matrix elements

Outlook

- ✓ Study CP-odd effects at colliders (e.g. Bernreuther et al, Tattersall et al)
- ✓ Include Higgs decay and differential distributions
- ✓ Extend analysis to Higgs-EW gauge bosons (in preparation)
- ✓ Compare linear v non-linear EFT realization

Backup

Bounds and scales

Use the neutron* EDM bound (**big uncertainty for some operators: that's why we are here !**)

Dekens, JdV JHEP '13

Dimensionless
couplings

	$M_T = 1 \text{ TeV}$	$M_T = 10 \text{ TeV}$
(M _T ²)d _{u, d} (M _T)	$\leq \{1.8, 1.8\} \cdot 10^{-3}$	$\leq \{2.1, 2.1\} \cdot 10^{-1}$
	$\leq \{1.9, 0.91\} \cdot 10^{-3}$	$\leq \{1.7, 0.94\} \cdot 10^{-1}$
	$\leq 5.6 \cdot 10^{-5}$	$\leq 7.0 \cdot 10^{-3}$
	$\leq 3.2 \cdot 10^{-5}$	$\leq 2.3 \cdot 10^{-3}$
	$\leq 3.3 \cdot 10^{-4}$	$\leq 2.4 \cdot 10^{-2}$
	$\leq 1.7 \cdot 10^{-4}$	$\leq 1.7 \cdot 10^{-2}$
	$\leq \{8.9, 8.9\} \cdot 10^{-5}$	$\leq \{7.9, 7.9\} \cdot 10^{-3}$
	$\leq 2.4 \cdot 10^{-3}$	$\leq 1.5 \cdot 10^{-1}$

* Hg EDM bound gives stronger limits for some operators (e.g. quark CEDM) but also suffers from larger theoretical uncertainty

Engel et al, PNPP '13

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	$\leq 2.4 \cdot 10^{-3}$	$\leq 1.5 \cdot 10^{-1}$

So 1 TeV seems ‘unnatural’ but note loop factors. For instance:

$$M_{CP}^2 \tilde{d}_q \sim \frac{\alpha_s}{4\pi} \sin \phi_{CP} \sim 10^{-2} \sin \phi_{CP} \quad \longrightarrow \quad \sin \phi_{CP} \leq 10^{-1}$$

The interpretation is model dependent

Bounds and scales

Use the neutron EDM bound (**big uncertainty for some operators:**
that's why we are here !)

Dekens, JdV JHEP '13

'electroweak suppressed operators'

	$M_T = 1 \text{ TeV}$	$M_T = 10 \text{ TeV}$
Dimensionless couplings	$(M_T^2)C_B(M_T)$	$\leq 8.1 \cdot 10^{-2}$
	$(M_T^2)C_W(M_T)$	$\leq 1.9 \cdot 10^{-2}$
	$(M_T^2)C_{WB}(M_T)$	$\leq 1.3 \cdot 10^{-2}$
	$(M_T^2)C_{dW}(M_T)$	≤ 0.11
	$(M_T^2)C_{Wu,d}(M_T)$	$\leq \{1.0, 0.84\} \cdot 10^{-2}$
	$(M_T^2)C_{Zu,d}(M_T)$	$\leq \{5.3, 2.8\} \cdot 10^{-2}$
		$\leq \{0.53, 0.45\}$
		$\leq \{2.7, 1.4\}$

First 4 operators better bound by eEDM

Lattice QCD to the rescue

- ❖ With QCD lattice input:

$$d_n = (2.7 \pm 1.2) \cdot 10^{-16} \overline{\theta} e \text{ cm} \quad \text{Shintani et al '12 '13}$$
$$d_p = -(2.1 \pm 1.2) \cdot 10^{-16} \overline{\theta} e \text{ cm}$$

$$d_n = (3.9 \pm 1.0) \cdot 10^{-16} \overline{\theta} e \text{ cm} \quad \text{Guo et al '15}$$

- ChPT extrapolation to **physical pion mass** and **infinite volume**

$$d_n = \overline{d}_0 - \overline{d}_1 - \frac{eg_A\overline{g}_0}{4\pi^2 F_\pi} \left(\ln \frac{m_\pi^2}{M_N^2} - \frac{\pi}{2} \frac{m_\pi}{M_N} \right) \quad \begin{array}{l} \text{O'Connell, Savage '06} \\ \text{Guo, Meißner, Akan '14} \end{array}$$

- Popular problem: see also Shindler et al '15
Alexandrou et al '15

Still not really clear though....

