Precision tools and models to narrow in on the 750 GeV diphoton resonance

Toby Opferkuch

Based on arXiv:1602.05581

in collaboration with

F. Staub, P. Athron, L. Basso, M. D. Goodsell, D. Harries, M. E. Krauss, K. Nickel, L. Ubaldi, A. Vicente and A. Voigt

ABHM Meeting Mainz



March 7, 2016

The purpose of this talk?

i) Emphasise a number of deficiencies in the diphoton literature

ii) Show how using **SARAH** framework can help to prevent these deficiencies

iii) Illustrate elements of a small complete example

DIPHOTON MODELS

Composite Models

- $\mathcal{O}(20)$ papers
- Naturally broad resonance



• $\mathcal{O}(10)$ papers



Perturbative Models

• $\mathcal{O}(200)$ papers

[All references in 1602.05581]

DIPHOTON MODELS

Composite Models

- $\mathcal{O}(20)$ papers
- Naturally broad resonance

Extra-dimensions

• $\mathcal{O}(10)$ papers



Perturbative Models

• $\mathcal{O}(200)$ papers

[All references in 1602.05581]

Decay width

Generic explanation involves loop-induced couplings to both photons and gluons 10 $b\overline{b} \rightarrow S$ ത്ത ↑ L₈/L₁₃ CMS 8 Sσ13 TeV/σ8 TeV $6 \uparrow L_8/L_{13}$ ATLAS ത്ത NP un [S. Knapen et al. 1512.04928] 500 1000 1500 2000 í٥ Important prediction of a model Mass of the resonance in GeV is the ratio $BR(S \to qq)/BR(S \to \gamma\gamma)$

[R. Franceschini et al 1512.04933]

DECAY WIDTH



Consider toy model containing:

Vector-like quarks Ψ (3,2,7/6) Singlet S (1,1,0)

$$\mathcal{L}_Y \supset (M_{F_1} + Y_{F_1}S) \overline{\Psi_L} \Psi_R + \text{h.c.}$$

[S. Knapen et al. 1512.04928]

DECAY WIDTH



Consider toy model containing:

Vector-like quarks Ψ (3, 2, 7/6) Singlet S (1, 1, 0)

$$\mathcal{L}_Y \supset (M_{F_1} + Y_{F_1}S) \overline{\Psi_L} \Psi_R + \text{h.c.}$$

[S. Knapen et al. 1512.04928]

CONCLUSION

- Small mismatch in LO result as $\alpha_{\rm em}(0) \neq \alpha_{\rm em}(M_S)$
- QCD corrections dominate over LO mismatch

Necessary decay rate depends on signal width

- Narrow width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-6}$
- Large width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-4}$

[R. Franceschini et al 1512.04933]

Necessary decay rate depends on signal width

- Narrow width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-6}$
- Large width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-4}$

[R. Franceschini et al 1512.04933]

How to increase the width:

- Fermions with large Yukawa couplings
- Fermions/scalars with large multiplicity and/or electric charge
- Scalars with large cubic couplings

Necessary decay rate depends on signal width

- Narrow width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-6}$
- Large width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-4}$

[R. Franceschini et al 1512.04933]

How to increase the width:

- I Fermions with large Yukawa couplings
- Fermions/scalars with large multiplicity and/or electric charge
- Scalars with large cubic couplings

Option 1

Pert. calculation \Rightarrow Yukawa must remain perturbative $\lesssim \sqrt{4\pi}$

(Not required in composite models)

Necessary decay rate depends on signal width

- Narrow width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-6}$
- Large width: $\Gamma(S \to \gamma \gamma)/M_S \simeq 10^{-4}$

[R. Franceschini et al 1512.04933]

How to increase the width:

- Fermions with large Yukawa couplings
- Fermions/scalars with large multiplicity and/or electric charge
- Scalars with large cubic couplings

Option 2

Must check for Landau poles Example: SM + N_k generations of k^{++}

[Kanemura et al. 1512.09048

Nomura et al. 1601.00386]

N_k	$\mu_{ m Landau}$
10	$2 \times 10^{13} {\rm TeV}$
100	$1.2 \times 10^5 \mathrm{TeV}$
1000	$3.8 { m TeV}$
6000	$2.7 { m TeV}$
9000	$2.6 { m TeV}$

Option 3

Alternative to fermions with large Yukawas \longrightarrow scalars with large κ

$$V \supset \kappa S|X|^2 + \frac{1}{2}M_S S^2 + M_X|X|^2 + \cdots$$

Stability of the EW vacuum:



MIXING WITH THE SM HIGGS

Toy model with CP-even singlet S and electrically charged scalars X

$$V = \frac{1}{2}M_S S^2 + M_X |X|^2 + \mu^2 |H|^2 + \kappa S |X|^2 + \kappa_H S |H|^2 + \lambda_S S^4 + \lambda_{SX} S^2 X^2 + \lambda_{HX} |H|^2 |X|^2 + \lambda |H|^4$$

AT TREE-LEVEL

- Many studies choose $\kappa_H = 0 \Longrightarrow$ no mixing
- Non-zero mixing:
 - \implies tree-level decays \gg loop decays

MIXING WITH THE SM HIGGS

Toy model with CP-even singlet S and electrically charged scalars X

$$V = \frac{1}{2}M_S S^2 + M_X |X|^2 + \mu^2 |H|^2 + \kappa S |X|^2 + \kappa_H S |H|^2 + \lambda_S S^4 + \lambda_{SX} S^2 X^2 + \lambda_{HX} |H|^2 |X|^2 + \lambda |H|^4$$

AT TREE-LEVEL

- Many studies choose $\kappa_H = 0 \Longrightarrow$ no mixing
- Non-zero mixing:
 - \implies tree-level decays \gg loop decays

LOOP-LEVEL



MIXING WITH THE SM HIGGS

Toy model with CP-even singlet S and electrically charged scalars X

$$V = \frac{1}{2}M_S S^2 + M_X |X|^2 + \mu^2 |H|^2 + \kappa S |X|^2 + \kappa_H S |H|^2 + \lambda_S S^4 + \lambda_{SX} S^2 X^2 + \lambda_{HX} |H|^2 |X|^2 + \lambda |H|^4$$

AT TREE-LEVEL

- Many studies choose $\kappa_H = 0 \implies$ no mixing
- Non-zero mixing:
 ⇒ tree-level decays ≫ loop decays

Other considerations

- Similar arguments hold for vector-like fermions
- $\kappa_H \neq 0$ often required for pseudo-scalar masses
- Good solution: use CP to forbid unwanted tree-level decays (see later slide)

LOOP-LEVEL



Typical assumption

 $v_S = 0$ at all orders

Typical assumption

 $v_S = 0$ at all orders

TADPOLE EQUATIONS

$$\frac{\partial V^{(1L)}}{\partial v_S} = T^{(1L)} = T^{(T)} + \delta T = 0$$

Assuming

$$T^{(T)} = c_1 v_S + c_2 v_S^2 + c_3 v_S^3 = 0$$

Typical assumption

 $v_S = 0$ at all orders

TADPOLE EQUATIONS

$$\frac{\partial V^{(1L)}}{\partial v_S} = T^{(1L)} = T^{(T)} + \delta T = 0$$

Assuming

$$T^{(T)} = c_1 v_S + c_2 v_S^2 + c_3 v_S^3 = 0$$

ONE-LOOP RESULTS

$$\delta T = \begin{cases} \kappa A(M_X^2) & \text{scalar loop} \\ 2Y M_{\Psi} A(M_{\Psi}^2) & \text{fermion loop} \end{cases}$$
$$A(x^2) = \frac{1}{16\pi^2} x^2 \left[1 + \log\left(\frac{\mu^2}{x^2}\right) \right]$$



Result

If $M_{\Psi} \sim \kappa \sim M_X \sim \mathcal{O}(1 \,\text{TeV})$ then $\delta T \simeq 1 \,\text{TeV}^3/(16\pi^2 c_1)$

WHAT IS SARAH AND HOW DOES IT HELP?

WHAT IS SARAH AND HOW DOES IT HELP?



WHAT IS SARAH AND HOW DOES IT HELP?

Can consider a complete model without (erroneous) simplifying assumptions

What's new?

- diphoton and digluon effective vertices calulated
- SPheno then calculates decay width and production x-sec.
- Eff. vertices can be passed to MadGraph & CalcHep

DECAY WIDTH IMPLEMENTATION

$\Phi \to \gamma \gamma$

LO expressions for decay width implemented using $\alpha_{ew}(\mu = 0)$

NLO SM corrections implemented for three limits:

- m_Φ < m_f: corrections from heavy colour fermionic triplets
- m_Φ > 100m_f: analytic corrections in light quark limit

 $[\mathrm{M.\ Spira\ et\ al.\ hep-ph}/9504378]$

• Intermediate range: numerical values from HDECAY used

[Djouadi et al. hep-ph/9704448]

DECAY WIDTH IMPLEMENTATION

$\Phi \to \gamma \gamma$

LO expressions for decay width implemented using $\alpha_{ew}(\mu = 0)$

NLO SM corrections implemented for three limits:

- m_Φ < m_f: corrections from heavy colour fermionic triplets
- m_Φ > 100m_f: analytic corrections in light quark limit

[M. Spira et al. hep-ph/9504378]

• Intermediate range: numerical values from HDECAY used

[Djouadi et al. hep-ph/9704448]

$\Phi \to gg$

LO expressions for decay width implemented

 $N^{3}LO$ SM corrections implemented

[Baglio et al. 1312.4788

Kramer et al. hep-ph/961127

Chetyrkin et al. hep-ph/9705240, hep-ph/0512060

Schroder and Steinhauser hep-ph/0512058

Baikov and Chetyrkin hep-ph/0604194]

Only N²LO corrections for pseudo-scalar Φ



[LHC Higgs Cross Section Working Group Collaboration, J. R. Andersen et al. 1307.1347]



[LHC Higgs Cross Section Working Group Collaboration, J. R. Andersen et al. 1307.1347]



[LHC Higgs Cross Section Working Group Collaboration, J. R. Andersen et al. 1307.1347]



[LHC Higgs Cross Section Working Group Collaboration, J. R. Andersen et al. 1307.1347]

Model	Name			
model	Name			
Toy models with vector-like fermions				
CP-even singlet	SM+VL/CPevenS			
CP-odd singlet	SM+VL/CPoddS			
Complex singlet	SM+VL/complexS			
Models based o	n the SM gauge-group			
Portal dark matter	SM+VL/PortalDM			
Scalar octet	SM-S-Octet	$\mathbb{A}^{(1)}$		
SU(2) triplet quark model	SM+VL/TripletQuarks			
Single scalar leptoquark	LQ/ScalarLeptoquarks			
Two scalar leptoquarks	LQ/TwoScalarLeptoquarks	$\mathbb{A}^{(3)}$		
Georgi-Machacek model	Georgi-Machacek			
THDM w. colour triplet	THDM+VL/min-3			
THDM w. colour octet	THDM+VL/min-8			
THDM-I w. exotic fermions	THDM+VL/Type-I-VL			
THDM-II w. exotic fermions	THDM+VL/Type-II-VL			
THDM-I w. SM-like fermions	THDM+VL/Type-I-SM-like-VL			
THDM-II w. SM-like fermions	THDM+VL/Type-II-SM-like-VL			
THDM w. scalar septuplet	THDM/ScalarSeptuplet			

 $\underline{\wedge}^{(1)}$ conflict with limits from $S \to jj$, $\underline{\wedge}^{(3)}$ we disagree with their diphoton rate

Model	Name	
U(1) Exte	ensions	
Dark $U(1)'$	U1Ex/darkU1	
Hidden $U(1)$	U1Ex/hiddenU1	
Simple $U(1)$	U1Ex/simpleU1	
Scotogenic $U(1)$	U1Ex/scotoU1	$\wedge^{(2)}$
Unconventional $U(1)_{B-L}$	U1Ex/BL-VL	_
Sample of $U(1)'$	U1Ex/VLsample	
flavour-nonuniversal charges	U1Ex/nonUniversalU1	
Leptophobic $U(1)$	U1Ex/U1Leptophobic	$\wedge^{(1)}$
Z' mimicking a scalar resonance	U1Ex/trickingLY	_
Non-abelian gauge-group	extensions of the SM	
LR without bidoublets	LRmodels/LR-VL	$\mathbb{A}^{(2)}$
LR with $U(1)_L \times U(1)_R$	LRmodels/LRLR	$\overline{\mathbb{A}}^{(2)}$
LR with triplets	LRmodels/tripletLR	
Dark LR	LRmodels/darkLR	
331 model without exotic charges	331/v1	
331 model with exotic charges	331/v2	
Gauged THDM	GTHDM	
Gauged THDM	GTHDM	

 $\underline{\wedge}^{(1)}$ conflict with limits from $S \to jj$, $\underline{\wedge}^{(2)}$ inconsistencies in charge assignments

Model	Name	
Supersymn	netric models	
NMSSM with vectorlike top	NMSSM+VL/VLtop	$\mathbb{A}^{(1)}$
NMSSM with 5 's	NMSSM+VL/5plets	
NMSSM with 10 's	NMSSM+VL/10plets	
NMSSM with 5 's & 10 's	NMSSM+VL/10plets	
NMSSM with 5 's and R pV	NMSSM+VL/5plets+RpV	
Broken MRSSM	brokenMRSSM	
U(1)'-extended MSSM	MSSM+U1prime-VL	
E_6 with extra $U(1)$	E6MSSMalt	

 $\underline{\wedge}^{(1)}$ conflict with limits from $S \to jj$

MODEL DETAILS

Model features

- Gauge sector extended by $U(1)_X$
- Tree-level Higgs mass enhancement (non-decoupling *D*-terms)
- $\bullet~\mathrm{CP}\text{-}\mathrm{odd}$ scalar acts as $750\,\mathrm{GeV}$ resonance
- Can potentially accommodate broad resonance

Model details

Model features

- Gauge sector extended by $U(1)_X$
- Tree-level Higgs mass enhancement (non-decoupling *D*-terms)
- $\bullet~\mathrm{CP}\text{-}\mathrm{odd}$ scalar acts as $750\,\mathrm{GeV}$ resonance
- Can potentially accommodate broad resonance

$$W = -W_{\text{Yuk}} + Y_{\nu} \hat{\nu} \hat{l} \hat{H}_{u} + \hat{S}(\lambda_{e} \hat{E} \hat{E} + \lambda_{u} \hat{U} \hat{U})$$

+ $Y_{x} \hat{\nu} \hat{\eta} \hat{\nu} + (\mu + \lambda \hat{S}) \hat{H}_{u} \hat{H}_{d} + \hat{S}(\xi + \lambda_{X} \hat{\eta} \hat{\eta})$
+ $M_{S} \hat{S} \hat{S} + \frac{1}{3} \kappa \hat{S} \hat{S} \hat{S} + \tilde{M}_{E} \hat{e} \hat{E} + \tilde{M}_{U} \hat{u} \hat{U}$
+ $M_{e} \hat{E} \hat{E} + M_{u} \hat{U} \hat{U} + Y'_{e} \hat{E} \hat{l} \hat{H}_{d} + Y'_{u} \hat{U} \hat{q} \hat{H}_{u}$

\mathbf{SF}	Gen.	$(\mathcal{G}_{\mathrm{SM}}, U(1)_X)$
$\hat{\nu}$	3	$(1, 1, 0, -\frac{1}{2})$
\hat{U}	3	$(ar{3}, m{1}, -rac{2}{3}, -rac{1}{2})$
$\hat{\bar{U}}$	3	$(ar{3}, 1, rac{2}{3}, rac{1}{2})$
Ê	3	$(1, 1, 1, \frac{1}{2})$
$\hat{\bar{E}}$	3	$(1, 1, -1, -\frac{1}{2})$
$\hat{\eta}$	1	(1, 1, 0, -1)
$\hat{ar{\eta}}$	1	(1, 1, 0, 1)
\hat{S}	1	$({f 1},{f 1},0,0)$

[R. M. Capdevilla, A. Delgado, and A. Martin 1509.02472]

Model details

Model features

- Gauge sector extended by $U(1)_X$
- Tree-level Higgs mass enhancement (non-decoupling *D*-terms)
- CP-odd scalar acts as 750 GeV resonance
- Can potentially accommodate broad resonance

$$\begin{split} W &= -W_{\text{Yuk}} + Y_{\nu} \,\hat{\nu} \,\hat{l} \,\hat{H}_{u} + \hat{S}(\lambda_{e} \,\hat{E} \,\hat{E} + \lambda_{u} \,\hat{U} \,\hat{U}) \\ &+ \frac{Y_{x}}{Y_{x}} \,\hat{\nu} \,\hat{\eta} \,\hat{\nu} + (\mu + \lambda \hat{S}) \,\hat{H}_{u} \,\hat{H}_{d} + \hat{S}(\xi + \lambda_{X} \,\hat{\eta} \,\hat{\eta}) \\ &+ M_{S} \,\hat{S} \,\hat{S} + \frac{1}{3} \kappa \,\hat{S} \,\hat{S} \,\hat{S} + \tilde{M}_{E} \hat{e} \hat{E} + \tilde{M}_{U} \hat{u} \hat{U} \\ &+ M_{e} \,\hat{E} \,\hat{E} + M_{u} \,\hat{U} \,\hat{U} + Y_{e}^{\prime} \,\hat{E} \,\hat{l} \,\hat{H}_{d} + Y_{u}^{\prime} \,\hat{U} \,\hat{q} \,\hat{H}_{u} \end{split}$$

[R. M. Capdevilla, A. Delgado, and A. Martin 1509.02472]

SHORT ANALYSIS INCLUDES:

- Full tree-level mass spectrum
- RGEs and gauge kinetic mixing
- Two-loop Higgs mass corrections
- Two-loop corrections to the 750 GeV scalar
- Diphoton and digluon rates
- Full scalar BRs and singlet-doublet mixing
- Compatibility with SM Higgs measurements
- Width constraints from vacuum stability
- DM relic abundance
- Constraints from rare lepton flavour processes
- Z' mass limits

RESONANCE DECAY MODES

Illustrate effect of singlet-doublet mixing $\lambda \neq 0$:

- CP-even scalar mainly mixture $\eta \& \bar{\eta}$ with small singlet component
- CP-odd almost purely singlet



RESONANCE DECAY MODES



ATLAS results slightly prefer a large width $\sim 40\,{\rm GeV}$

Explained with inv. decays to:

- Neutralinos
- Heavy neutrinos
- Sneutrinos

ATLAS results slightly prefer a large width $\sim 40\,{\rm GeV}$

Explained with inv. decays to:

- Neutralinos
- Heavy neutrinos
- Sneutrinos

IN THIS MODEL

Small sneutrino masses and splitting between Re & Im can be achieved BR to sneutrinos scales with Y_x

ATLAS results slightly prefer a large width $\sim 40\,{\rm GeV}$

Explained with inv. decays to:

- Neutralinos
- Heavy neutrinos
- Sneutrinos

IN THIS MODEL

Small sneutrino masses and splitting between Re & Im can be achieved BR to sneutrinos scales with Y_x

QUESTIONS?

• How large can Γ_{tot} varying Y_x ?

• How large can Y_x be before the vacuum becomes unstable?

ATLAS results slightly prefer a large width $\sim 40\,{\rm GeV}$

Explained with inv. decays to:

- Neutralinos
- Heavy neutrinos
- Sneutrinos

Using Vevacious:



IN THIS MODEL

Small sneutrino masses and splitting between Re & Im can be achieved BR to sneutrinos scales with Y_x

QUESTIONS?

- How large can Γ_{tot} varying Y_x ?
- How large can Y_x be before the vacuum becomes unstable?



SUMMARY

SARAH framework allows easy analysis of complete models

Reduces necessity to make (extreme) simplifying assumptions

Introduced perturbative model where large width is feasible

BACKUP SLIDES

TREE-LEVEL VS. ONE-LOOP

CAUTION

Tree-level enforced relations (w/o symmetry arguments) **do not** hold at the loop-level



TREE-LEVEL VS. ONE-LOOP

CAUTION

Tree-level enforced relations (w/o symmetry arguments) **do not** hold at the loop-level

$\Gamma_{XX}/\Gamma_{\gamma\gamma}$
0.6
6
6
6
10
20
20
300
500
1300
100



Constraints from $8 \,\mathrm{TeV}$ run



PARAMETER VALUES

Mixing and decay width plots:

$$m_{\rm SUSY} = 1.5 \text{ TeV}, M_{\lambda} = 1 \text{ TeV}, \tan \beta = 20, \tan \beta_x = 1, g_X = 0.5, M_{Z'} = 3 \text{ TeV}, \\ \mu = 1 \text{ TeV}, B_{\mu} = (1 \text{ TeV})^2, v_S = 0.5 \text{ TeV}, M_S = -0.1 \text{ TeV}, B_S = 3.895 \text{ TeV}^2, \\ \lambda_X = -0.2, A_X = 1 \text{ TeV}, \lambda_E = \lambda_U = 1, M_E = 0.4 \text{ TeV}, M_U = 1 \text{ TeV}, m_{\bar{\eta}} = 2 \text{ TeV}.$$

For vacuum stability:

$$\begin{split} m_{\rm SUSY} &= 2.5~{\rm TeV}\,, \tan\beta = 10\,, \tan\beta_x = 1\,, g_X = 0.5\,, M_{Z'} = 2.5~{\rm TeV}, m_{\bar{\eta}} = 1~{\rm TeV}, \\ v_S &= 0.5~{\rm TeV}\,, B_S = 755000~{\rm GeV}^2\,, \lambda_X = -0.4\,, A_X = 0.4~{\rm TeV}. \end{split}$$